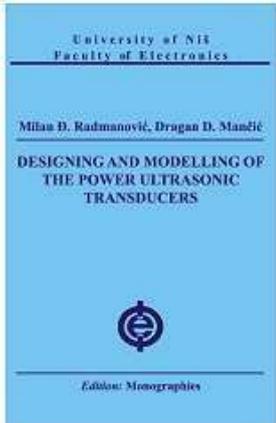


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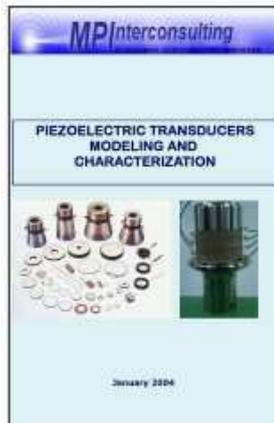
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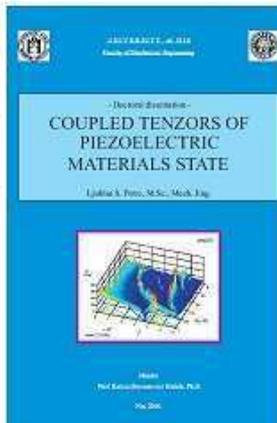
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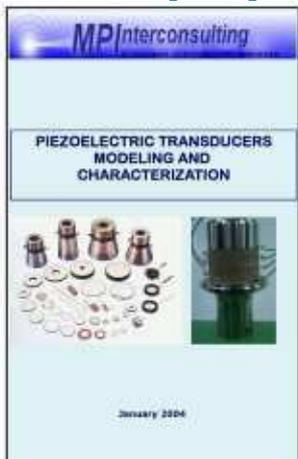


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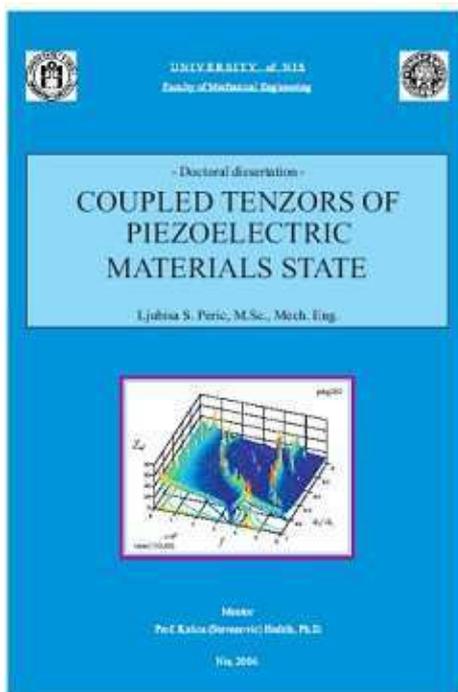
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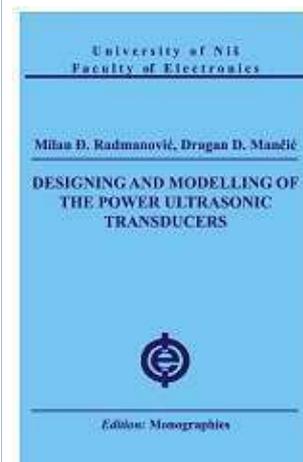
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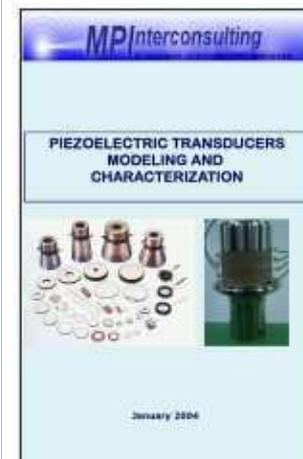
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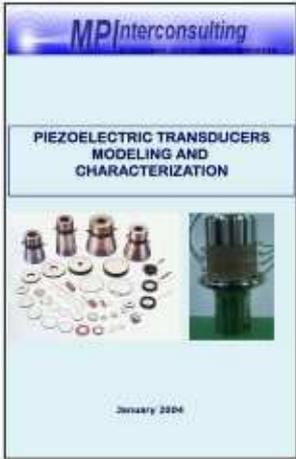
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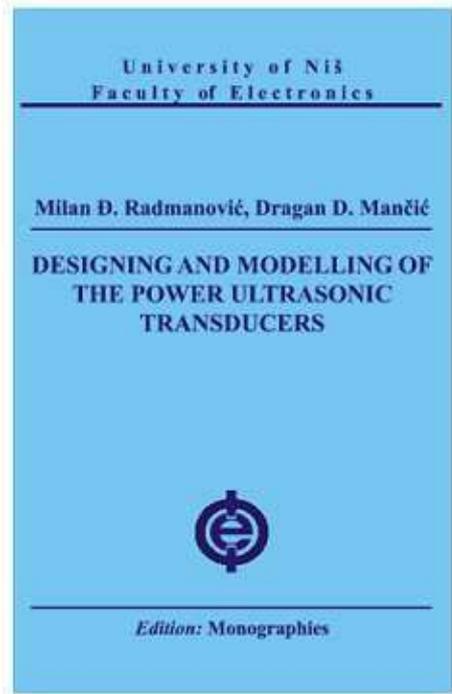
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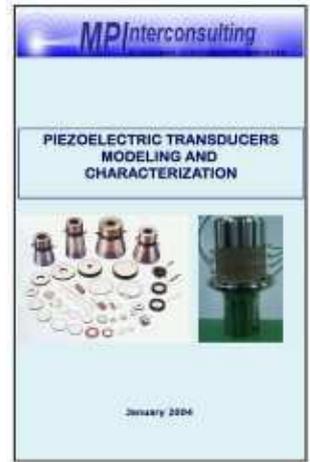
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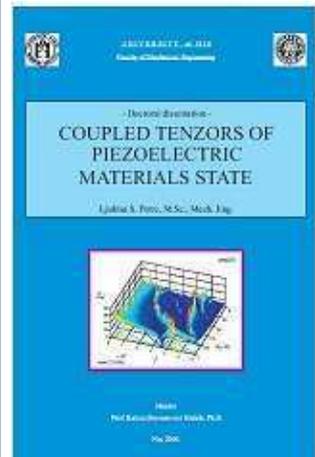


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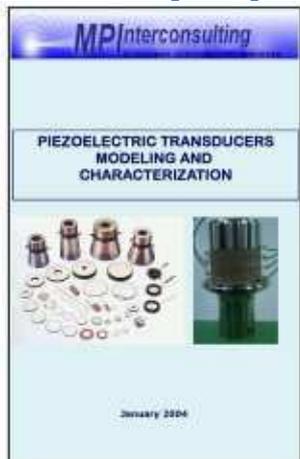


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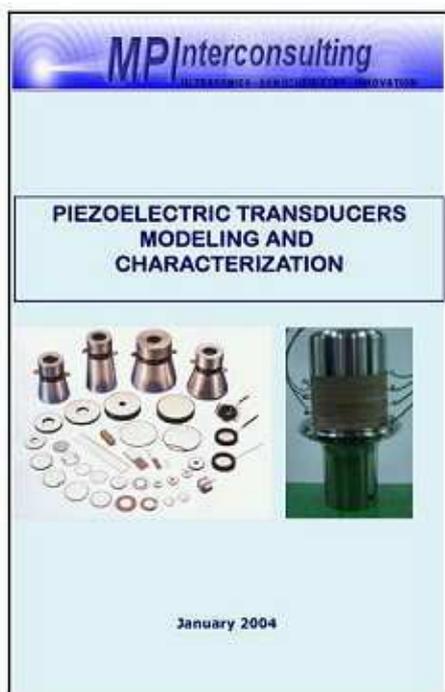
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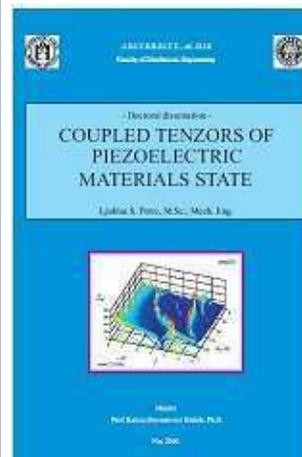
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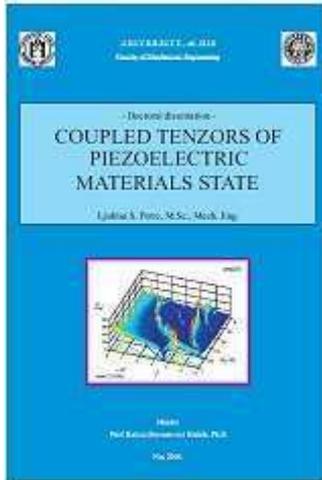
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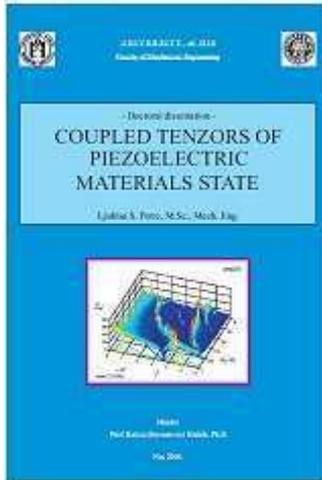
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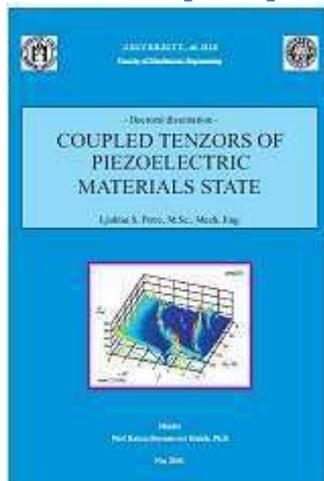
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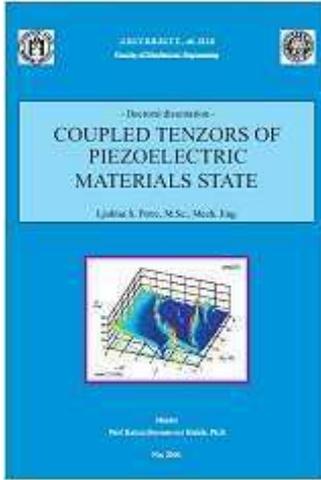
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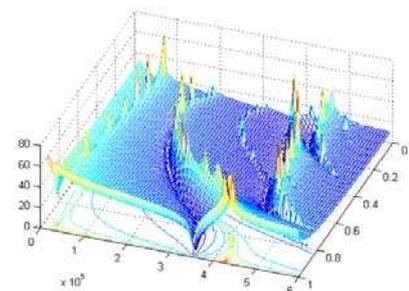
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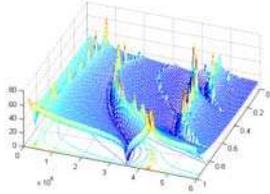
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BOOK: Coupled Tensors of Piezoelectric
Materials State and Applications.

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Universal Formula for Electric Impedance of Piezoelectric Ceramic Z_p

I have derived an **Universal Formula for Simulation of Electric Impedance Model** $Z_p = V/I$ [Ω] of piezoelectric ceramic. Formula stands for all shapes of piezoelectric bodies: circular plate, circular-ring plate, rectangular plate, cylinder, circular-ring cylinder, circular-ring sectional cylinder, sphere, etc. Piezoceramic bodies may have n -th number of surfaces loaded by equal or different external loads: F_j^D or F_j^E (v_j^D or v_j^E - motion velocities of contour surfaces loaded by F_j^D or F_j^E) as interaction with outer medium, i.e., external mechanical impedances: Z_j^D or Z_j^E .

Formula enables precise determination of resonant frequencies of numerous radial, transversal, longitudinal, and lateral modes of oscillation even at design stage of piezoelectric transducers, and before manufacturing of piezoceramic elements and experimental measuring. Application field of piezo-sensors and actuators is very wide, and of special interest for **military industry** and **space research**. It may be also used for significant improvement of existing methods of modeling FEM, BEM, etc.

The formula is based on two constitutive system of equations:

Isothermal func. of internal energy $U(S_{ij}, D_i)$:	Isothermal func. of electric potential $H(S_{ij}, E_i)$:
$T_{rr}^D = c_{11}^D S_{rr} + c_{12}^D S_{\theta\theta} + c_{13}^D S_{zz} - h_{31} D_z,$	$T_{rr}^E = c_{11}^E S_{rr} + c_{12}^E S_{\theta\theta} + c_{13}^E S_{zz} - e_{31} E_z,$
$T_{\theta\theta}^D = c_{12}^D S_{rr} + c_{11}^D S_{\theta\theta} + c_{13}^D S_{zz} - h_{31} D_z,$	$T_{\theta\theta}^E = c_{12}^E S_{rr} + c_{11}^E S_{\theta\theta} + c_{13}^E S_{zz} - e_{31} E_z,$
$T_{zz}^D = c_{13}^D (S_{rr} + S_{\theta\theta}) + c_{33}^D S_{zz} - h_{33} D_z,$	$T_{zz}^E = c_{13}^E (S_{rr} + S_{\theta\theta}) + c_{33}^E S_{zz} - e_{33} E_z,$
$E_z^D = -h_{31} (S_{rr} + S_{\theta\theta}) - h_{33} S_{zz} + D_z / \epsilon_{33}^S,$	$D_z^E = e_{31} (S_{rr} + S_{\theta\theta}) + e_{33} S_{zz} + \epsilon_{33}^S E_z.$

Formula of electric impedance Z_p , for piezoelectric body with n surfaces is a function of numerous parameters.

Sum of all forms of energy in a conservative system is constant (First Law of Thermodynamics, or Law of Conservation of Energy):

$$\sum_{j=1}^{n-1} F_j^D v_j^D + V^D I = \sum_{j=1}^{n-1} F_j^E v_j^E + I^E V, \quad n = 2, 3, 4, \dots, k.$$

Then follows:

$$Z_p = \frac{V}{I} = \sqrt{\frac{Z_{mm}^D + \sum_{i=1}^{n-1} Z_{in}^D \bar{v}_i^D + \sum_{j=1}^{n-1} Z_{nj}^D \bar{v}_j^D + \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} Z_{ij}^D \bar{v}_i^D \bar{v}_j^D}{Z_{mm}^E + \sum_{i=1}^{n-1} Z_{in}^E \bar{v}_i^E + \sum_{j=1}^{n-1} Z_{nj}^E \bar{v}_j^E + \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} Z_{ij}^E \bar{v}_j^E \bar{v}_i^E}}, \quad n = 2, 3, 4, \dots, k.$$

$Z_{ij}^D, Z_{nj}^D, Z_{mm}^D, Z_{ij}^E, Z_{nj}^E$ and Z_{mm}^E - internal mechanical and electric impedances, transfer functions of system (black box).

Comparison of derived formula by computer simulation using software package MATLAB, with analogous characteristic obtained by experimental measuring on Automatic Network Analyzer HP4194A, shows results at least 50% better than all known results published until now in scientific literature available to me.

Formula is applicable on: 1D, 2D, and 3D oscillating models of piezoelectric ceramic bodies, whose number of surfaces $S_n [m^2]$ may extend to infinity ($n \rightarrow \infty$). Also, it can join different assumed oscillation models (any combination of: 1D, 2D, and 3D model) of the same piezoceramic specimen loaded under equal conditions in the same period of time.

1 – Case: $n = 2$ - **1D model** (transversal oscillation of the: circular plate, circular-ring plate, rectangular plate etc; or longitudinal oscillation of a: free beam, cantilever, cylinder, circular-ring cylinder, circular-ring sectional cylinder, etc.):

$$F_1^D v_1^D + V^D I = F_1^E v_1^E + I^E V,$$

$$Z_p = \frac{V}{I} = \sqrt{\frac{(z_{11}^D \bar{v}_1^D + z_{12}^D) \bar{v}_1^D + z_{21}^D \bar{v}_1^D + z_{22}^D}{(z_{11}^E \bar{v}_1^E + z_{12}^E) \bar{v}_1^E + z_{21}^E \bar{v}_1^E + z_{22}^E}}.$$

2 – Case: $n = 3$ - **2D model** (radial oscillation of the: circular plate, circular-ring plate, rectangular plate, free beam, cantilever, cylinder, circular-ring cylinder, circular-ring sectional cylinder, etc.):

$$F_1^D v_1^D + F_2^D v_2^D + V^D I = F_1^E v_1^E + F_2^E v_2^E + I^E V,$$

$$Z_p = \frac{V}{I} = \sqrt{\frac{(z_{11}^D \bar{v}_1^D + z_{12}^D \bar{v}_2^D + z_{13}^D) \bar{v}_1^D + (z_{21}^D \bar{v}_1^D + z_{22}^D \bar{v}_2^D + z_{23}^D) \bar{v}_2^D + z_{31}^D \bar{v}_1^D + z_{32}^D \bar{v}_2^D + z_{33}^D}{(z_{11}^E \bar{v}_1^E + z_{12}^E \bar{v}_2^E + z_{13}^E) \bar{v}_1^E + (z_{21}^E \bar{v}_1^E + z_{22}^E \bar{v}_2^E + z_{23}^E) \bar{v}_2^E + z_{31}^E \bar{v}_1^E + z_{32}^E \bar{v}_2^E + z_{33}^E}}.$$

3 – Case: $n = 4$ - **3D model** (3D oscillation of the: circular plate, circular-ring plate, rectangular plate, free beam, cantilever, cylinder, circular-ring cylinder, circular-ring sectional cylinder, etc.):

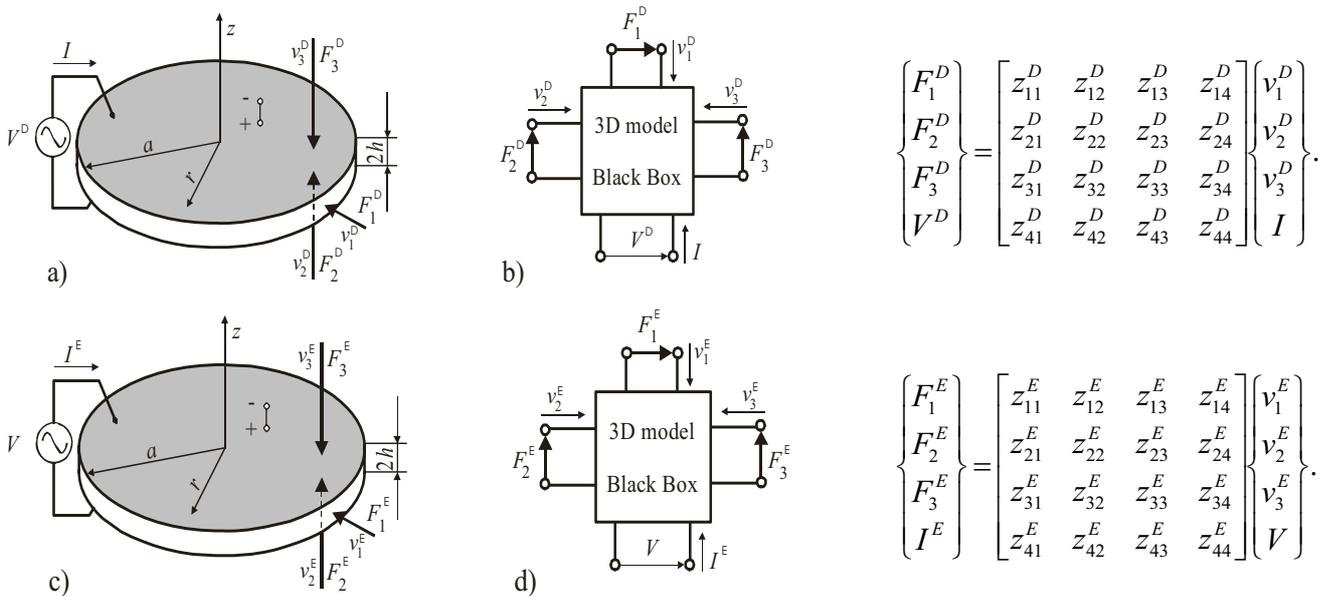
$$F_1^D v_1^D + F_2^D v_2^D + F_3^D v_3^D + V^D I = F_1^E v_1^E + F_2^E v_2^E + F_3^E v_3^E + I^E V,$$

$$Z_p = \frac{V}{I} = \sqrt{\frac{(z_{11}^D \bar{v}_1^D + z_{12}^D \bar{v}_2^D + z_{13}^D \bar{v}_3^D + z_{14}^D) \bar{v}_1^D + (z_{21}^D \bar{v}_1^D + z_{22}^D \bar{v}_2^D + z_{23}^D \bar{v}_3^D + z_{24}^D) \bar{v}_2^D + (z_{31}^D \bar{v}_1^D + z_{32}^D \bar{v}_2^D + z_{33}^D \bar{v}_3^D + z_{34}^D) \bar{v}_3^D + z_{41}^D \bar{v}_1^D + z_{42}^D \bar{v}_2^D + z_{43}^D \bar{v}_3^D + z_{44}^D}{(z_{11}^E \bar{v}_1^E + z_{12}^E \bar{v}_2^E + z_{13}^E \bar{v}_3^E + z_{14}^E) \bar{v}_1^E + (z_{21}^E \bar{v}_1^E + z_{22}^E \bar{v}_2^E + z_{23}^E \bar{v}_3^E + z_{24}^E) \bar{v}_2^E + (z_{31}^E \bar{v}_1^E + z_{32}^E \bar{v}_2^E + z_{33}^E \bar{v}_3^E + z_{34}^E) \bar{v}_3^E + z_{41}^E \bar{v}_1^E + z_{42}^E \bar{v}_2^E + z_{43}^E \bar{v}_3^E + z_{44}^E}}.$$

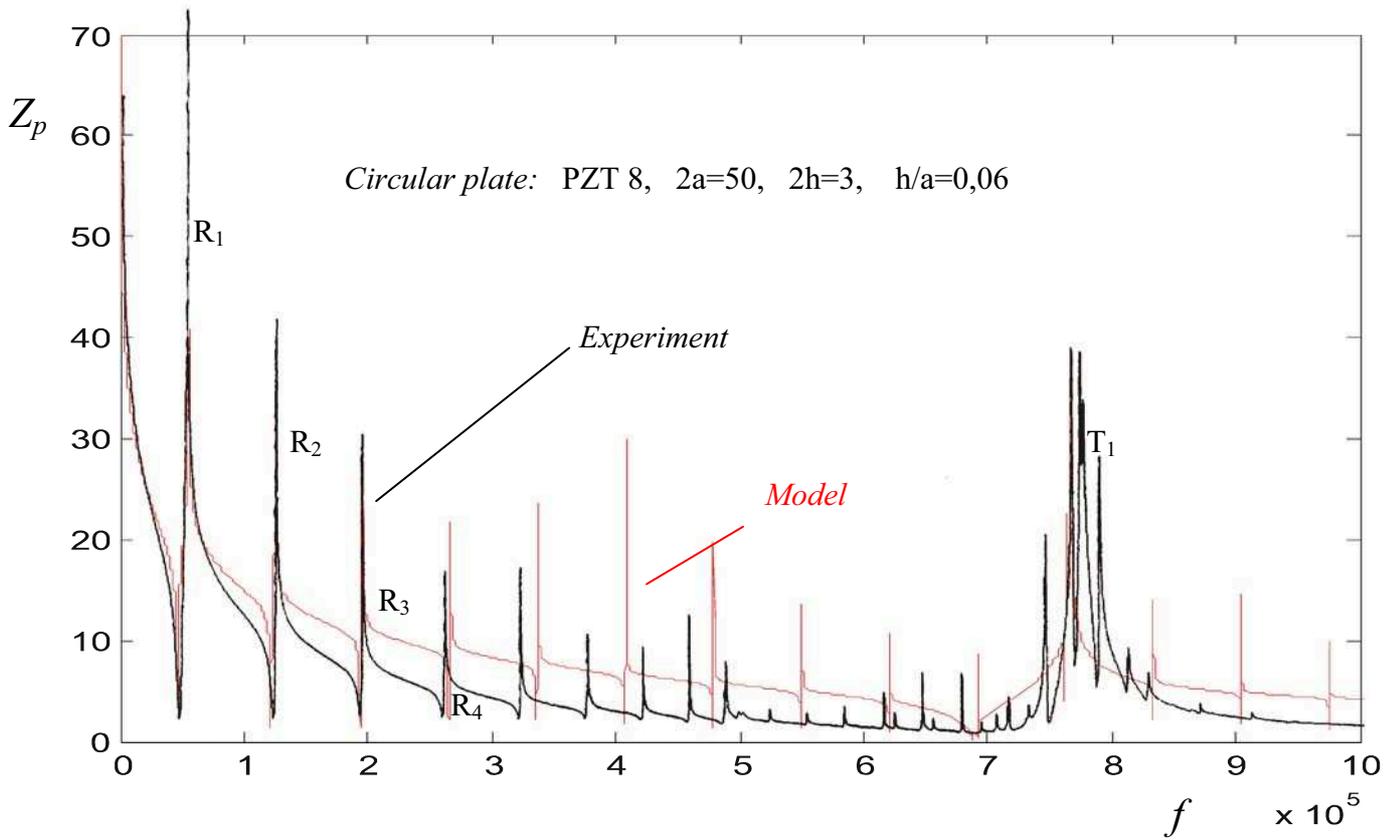
For PZT8 circular plate with dimensions: $2a_2 = 50$ [mm], $2h = 3$ [mm], $\rho = 7600$ [$\frac{kg}{m^3}$], whose coefficients are:

$$c_{11}^E = 13,7 \cdot 10^{10} [N/m^2], \quad c_{12}^E = 6,97 \cdot 10^{10}, \quad c_{13}^E = 7,16 \cdot 10^{10}, \quad c_{33}^E = 12,4 \cdot 10^{10}, \quad h_{31} = -7,8 \cdot 10^8 [V/m], \quad h_{33} = 26,9 \cdot 10^8 [V/m],$$

$$c_{11}^D = 14 \cdot 10^{10} [N/m^2], \quad c_{12}^D = 7,28 \cdot 10^{10}, \quad c_{13}^D = 6,08 \cdot 10^{10}, \quad c_{33}^D = 16,1 \cdot 10^{10}, \quad e_{31} = -4 [C/m^2], \quad e_{33} = 13,8, \quad \epsilon_{33}^S / \epsilon_0 = 582,$$



Formula provides following simulated characteristic of electric impedance model, which is compared with analogous, obtained by measuring on Automatic Network Analyzer HP4194A:



Knowing values of resonant frequencies is an initial condition during design of piezoceramic elements. From the picture above one may see that by this formula one may determine several exact positions of radial resonant modes R_1 , R_2 , R_3 and R_4 , and transversal mode T_1 , which are the most frequently used in technical practice and application.

4 – Case: $n = 5$ - 3D model (rectangular plate, free beam, etc.):

$$F_1^D v_1^D + F_2^D v_2^D + F_3^D v_3^D + F_4^D v_4^D + V^D I = F_1^E v_1^E + F_2^E v_2^E + F_3^E v_3^E + F_4^E v_4^E + I^E V,$$

$$Z_p = \frac{V}{I} = \sqrt{\frac{(z_{11}^{D-D} + z_{12}^{D-D} + z_{13}^{D-D} + z_{14}^{D-D} + z_{15}^D) \bar{v}_1^D + \dots + (z_{41}^{D-D} + z_{42}^{D-D} + z_{43}^{D-D} + z_{44}^{D-D} + z_{45}^D) \bar{v}_4^D + z_{51}^{D-D} + z_{52}^{D-D} + z_{53}^{D-D} + z_{54}^{D-D} + z_{55}^D}{(z_{11}^{E-E} + z_{12}^{E-E} + z_{13}^{E-E} + z_{14}^{E-E} + z_{15}^E) \bar{v}_1^E + \dots + (z_{41}^{E-E} + z_{42}^{E-E} + z_{43}^{E-E} + z_{44}^{E-E} + z_{45}^E) \bar{v}_4^E + z_{51}^{E-E} + z_{52}^{E-E} + z_{53}^{E-E} + z_{54}^{E-E} + z_{55}^E}}$$

5 – Case: $n = 6$ - 3D model (rectangular plate with a circular or rectangular hole, etc.):

$$F_1^D v_1^D + F_2^D v_2^D + F_3^D v_3^D + F_4^D v_4^D + F_5^D v_5^D + V^D I = F_1^E v_1^E + F_2^E v_2^E + F_3^E v_3^E + F_4^E v_4^E + F_5^E v_5^E + I^E V,$$

$$Z_p = \frac{V}{I} = \sqrt{\frac{(z_{11}^{D-D} + z_{12}^{D-D} + z_{13}^{D-D} + z_{14}^{D-D} + z_{15}^{D-D} + z_{16}^D) \bar{v}_1^D + \dots + (z_{51}^{D-D} + z_{52}^{D-D} + z_{53}^{D-D} + z_{54}^{D-D} + z_{55}^{D-D} + z_{56}^D) \bar{v}_6^D + z_{61}^{D-D} + z_{62}^{D-D} + z_{63}^{D-D} + z_{64}^{D-D} + z_{65}^{D-D} + z_{66}^D}{(z_{11}^{E-E} + z_{12}^{E-E} + z_{13}^{E-E} + z_{14}^{E-E} + z_{15}^{E-E} + z_{16}^E) \bar{v}_1^E + \dots + (z_{51}^{E-E} + z_{52}^{E-E} + z_{53}^{E-E} + z_{54}^{E-E} + z_{55}^{E-E} + z_{56}^E) \bar{v}_6^E + z_{61}^{E-E} + z_{62}^{E-E} + z_{63}^{E-E} + z_{64}^{E-E} + z_{65}^{E-E} + z_{66}^E}}$$

..., etc.

8.5.4. Second Approach to 3D Problem of Oscillations of Circular-ring Plate

In this part of the paper are considered 3D spatial oscillations of circular-ring plate with electrode coatings and transversal polarization along axis z Figure 8.121 – a. 3D model by which can be described spatial oscillation of piezoceramic circular-ring plate, may be presented using the black box analogy, Figure 8.121 – b.

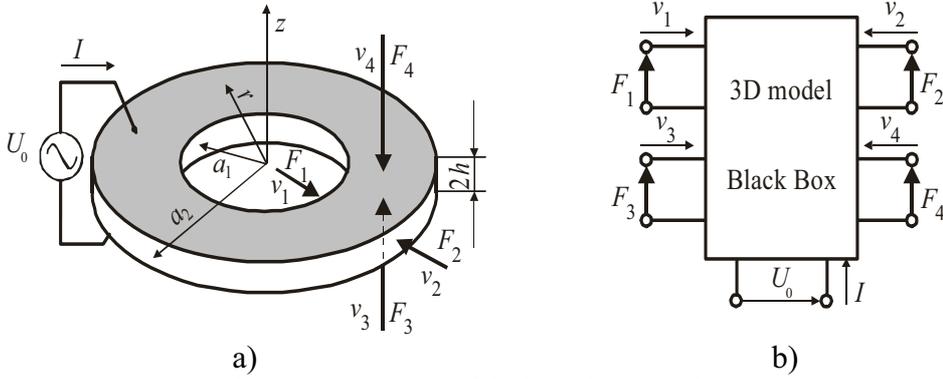


Figure 8.121. – 3D model of circular-ring plate

It may be seen that *circular-ring plate* may be observed as *black box* device with five variable parameters (five input-output values). First group of values (usually used as input – leading of electric energy on plate coatings from electric generator of alternate voltage) is *electric voltage* U_0 and *electric current strength* I . Remained four groups of values conform to cylindric surfaces - $r = a_1$ (*surface force* F_1 and *velocity* v_1), $r = a_2$ (*surface force* F_2 and *velocity* v_2), and plane mutually opposite surfaces $z = -h$ (*surface force* F_3 and *velocity* v_3), and $z = h$ (*surface force* F_4 and *velocity* v_4).

Oscillations of circular-ring plate are excited by leading of alternate difference of electric potential $2U_0 e^{i\omega t}$ on surfaces $z = \pm h$. *Plate oscillations*, i.e., *vibration of particles*, have *radial-transversal character of motion*, that is $\vec{s} = u(r, t)\vec{r}_0 + w(z, t)\vec{k}$.

Equations of piezoelectric effect, which are used in this analysis and procedure of derivation is identical with procedure in item 8.5.2*, expressions (8.271) ÷ (8.275).

Partial differential equations of oscillation of circular-ring plate:

$$c_{11}^D \left(\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} - \frac{u}{r^2} \right) = \rho \frac{\partial^2 u}{\partial t^2}, \quad (8.321)$$

$$c_{33}^D \frac{\partial^2 w}{\partial z^2} = \rho \frac{\partial^2 w}{\partial t^2}$$

* 8.5.1. 3D model of oscillations of circular plate

8.5.5.1. Diagrams of Spatial States

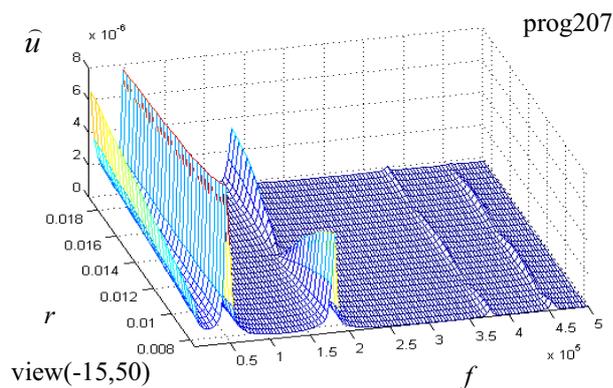


Figure 8.122. – Componential displacement $\hat{u} = \hat{u}(f, r)$

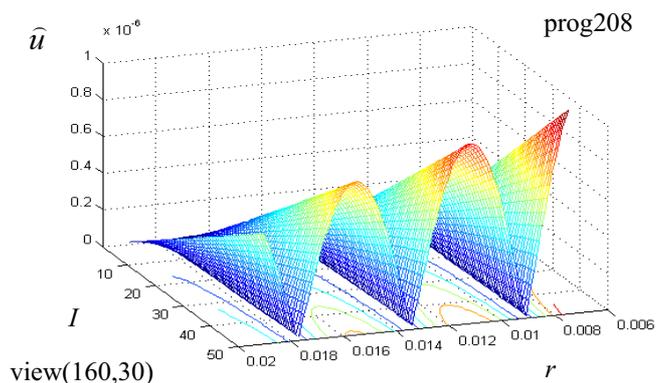


Figure 8.123. – Componential displacement $\hat{u} = \hat{u}(r, I)$

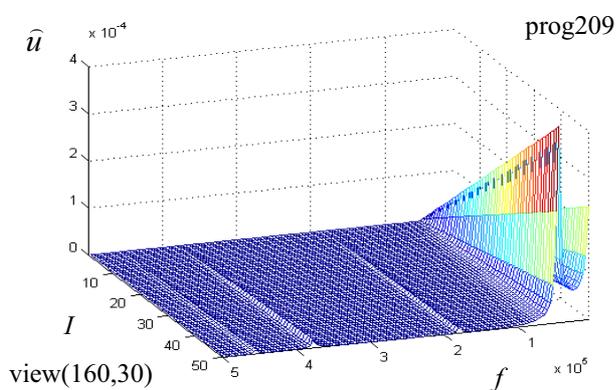


Figure 8.124. – Componential displacement $\hat{u} = \hat{u}(f, I)$

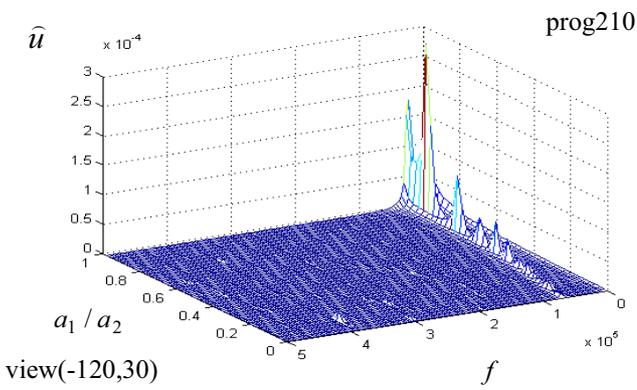


Fig. 8.125. – Component. Displacement $\hat{u} = \hat{u}(f, a_1 / a_2)$

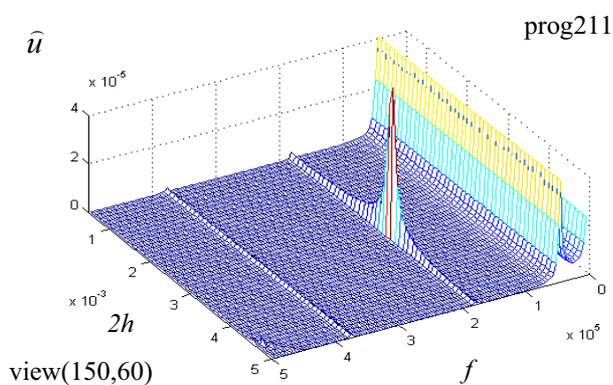


Figure 8.126. – Componential displacement $\hat{u} = \hat{u}(f, 2h)$

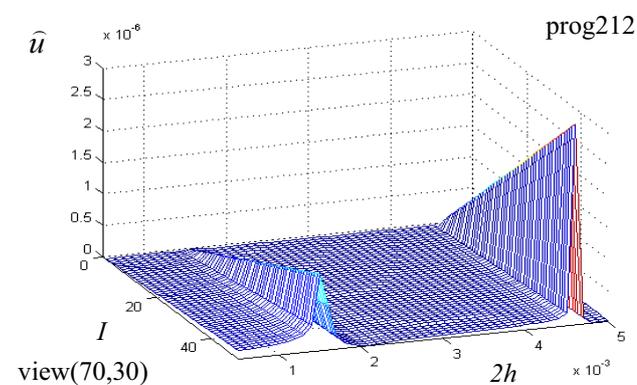
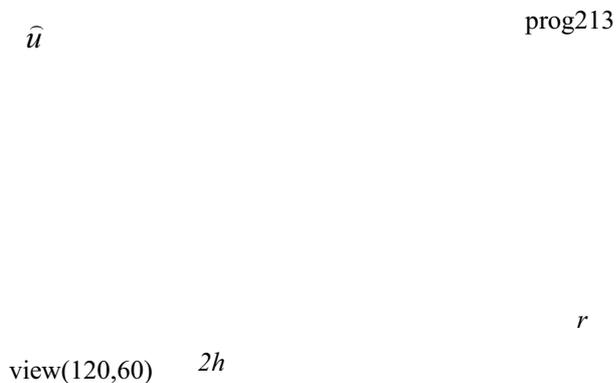


Figure 8.127. – Componential displacement $\hat{u} = \hat{u}(I, 2h)$



prog213

I

r

a_1 / a_2

$2h$

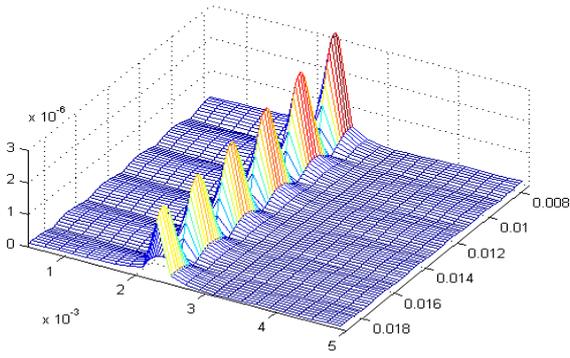


Figure 8.128. – Componential displacement $\hat{u} = \hat{u}(r, 2h)$

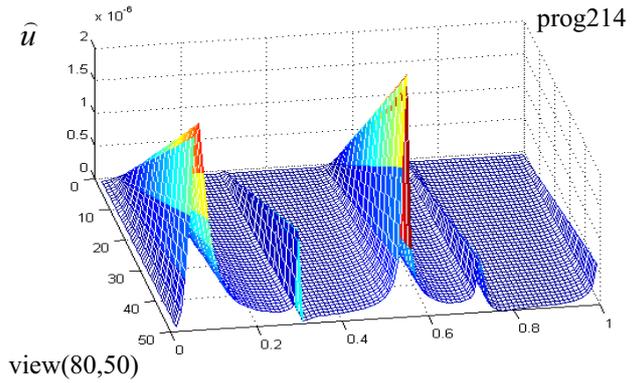


Fig. 8.129. – Component. Displacement $\hat{u} = \hat{u}(I, a_1 / a_2)$

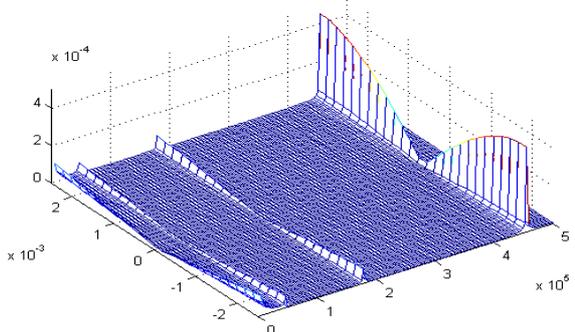


Figure 8.130. – Componential displacement $\hat{w} = \hat{w}(f, z)$

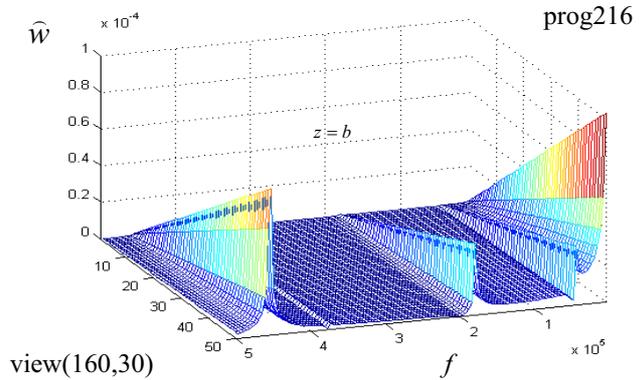
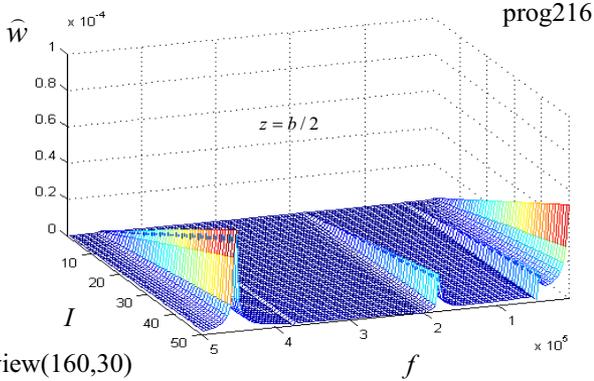
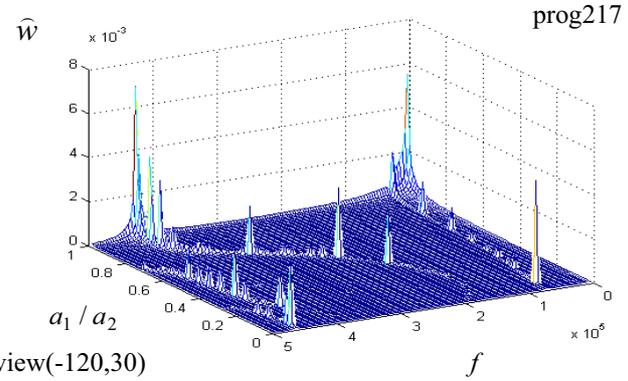


Figure 8.131. – Componential displacement $\hat{w} = \hat{w}(f, I)$



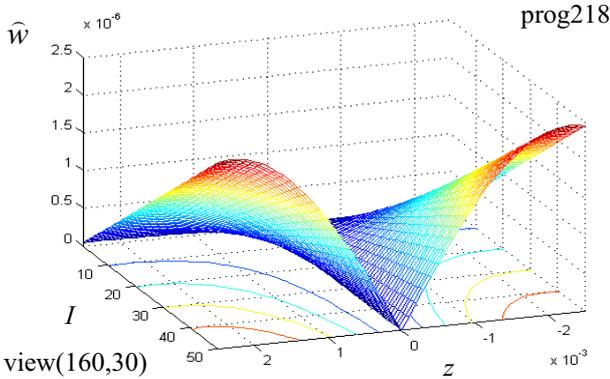
view(160,30)



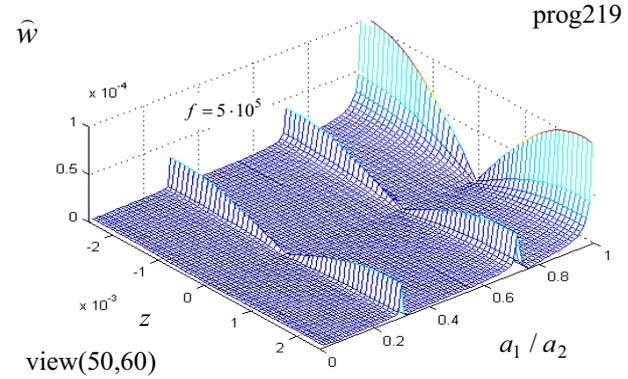
view(-120,30)

Figure 8.132. – Componential displacement $\hat{w} = \hat{w}(f, I)$

Fig. 8.133. – Component. Displacement $\hat{w} = \hat{w}(f, a_1 / a_2)$



view(160,30)



view(50,60)

Figure 8.134. – Componential displacement $\hat{w} = \hat{w}(z, I)$

Fig. 8.135. – Component. Displacement $\hat{w} = \hat{w}(z, a_1 / a_2)$

\hat{w} prog219

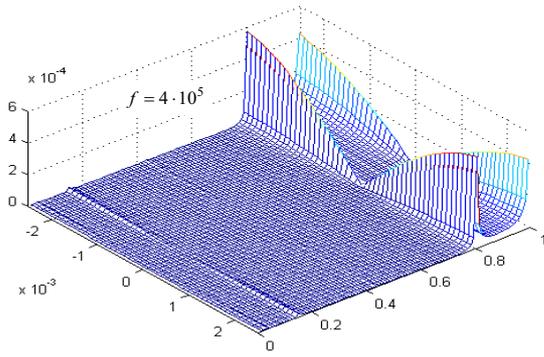


Fig. 8.136. – Component. Displacement $\hat{w} = \hat{w}(z, a_1 / a_2)$

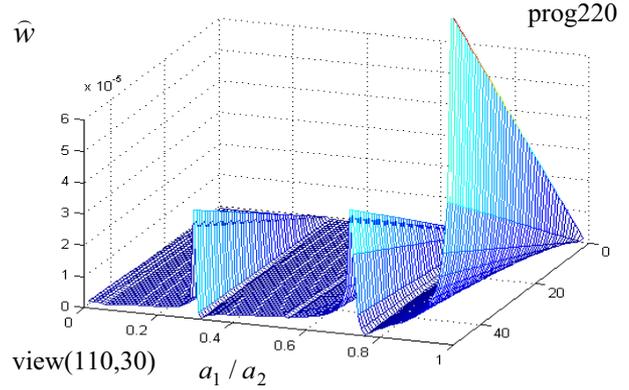


Fig. 8.137. – Component. Displacement $\hat{w} = \hat{w}(I, a_1 / a_2)$

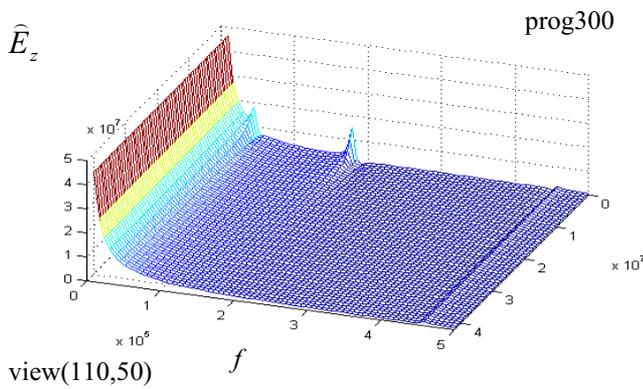


Fig. 8.274. – Electric field $\hat{E}_z = \hat{E}_z(f, Z_1 / Z_2)$

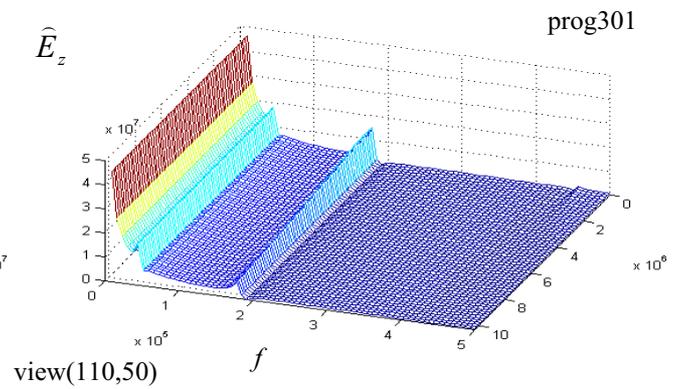


Fig. 8.275. – Electric field $\hat{E}_z = \hat{E}_z(f, Z_3 / Z_4)$

In order to show the possibility of analysis of piezoceramic circular-ring elements of concrete dimensions, a numerical analysis of the proposed model $Z_{ul} = Z_{ul}(f)$ is performed, using software package MATLAB. Concretely, an input electric impedance is determined for PZT8 piezoceramic circular-ring plate (Figure 8.276, programme 200 – Appendix III), whose tensors of material coefficients are given in Appendix II, with dimensions $2a_2 = 38[mm]$, $2a_1 = 15[mm]$, and $2h = 5[mm]$. It is assumed that all contour surfaces of circular-ring plate oscillate freely in air without additional external loading.

Comparison of obtained input electric impedance from Figure 8.276. is performed with analogous characteristic obtained by standard *Mason's* transversal (crossing) one-dimensional model [306] (Figure 8.277). One may notice that for the *first transversal mode* (T_1) exists satisfactory coincidence of both models. Small deviations that exist are result of coupled effect of action between transversal and radial mode of oscillation at adopted 3D model. Further, one may notice that one-dimensional *Mason's model* does not encompass radial resonant modes (R_1, R_2, R_3, R_4), which presents the proposed 3D model.

On Figure 8.278. is presented *comparative characteristic* of input electric impedance of adopted 3D model with characteristic of radial two-dimensional model [222]. One may notice that for the first two radial resonant modes of oscillation (R_1, R_2), exists good agreement, while for the rest radial resonant modes is characteristic deviation due to the great effect of transversal resonant mode at proposed 3D model. From Figure 8.278. one may notice that *two-dimensional radial model* cannot present *transversal mode of oscillation* (T_1).

On Figure 8.279. is presented comparison of electric impedance characteristics of adopted 3D model with 3D model proposed by *Brissaud* [53, 54]. One may notice that *Brissaud's model* contains certain limits and faults. Good coincidence is only at the first radial resonant mode (R_1), while at remained radial resonant modes deviations are significant. Also is illogical phenomenon at radial mode R_4 (*Brissaud's model*), that *characteristic of electric impedance* gets maximum first, and then minimum (region marked with arrow), which is not feasible and real. Beside the quoted faults of *Brissaud's* 3D model, one more is noticed, that it cannot encompass effect of external mechanical loads on boundary contour surfaces of circular-ring piezoceramic plate. Cited facts confirm great advantage of the proposed 3D model,* which is detailedly analyzed and derived in previous item (8.5.4).

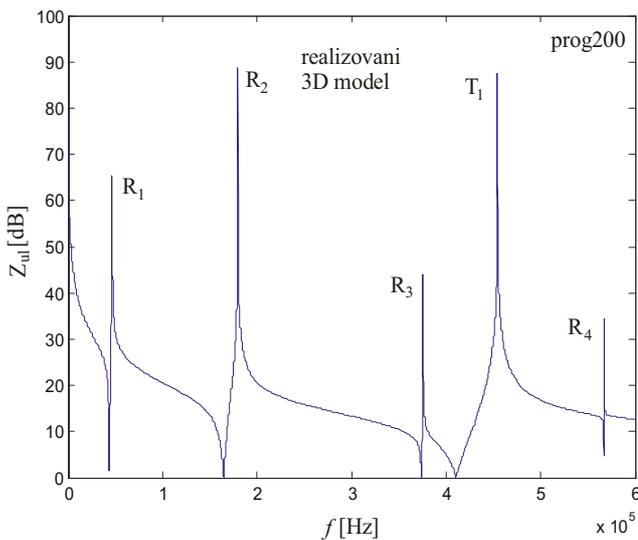


Figure 8.276. - Electric impedance $Z_{ul} = Z_{ul}(f)$

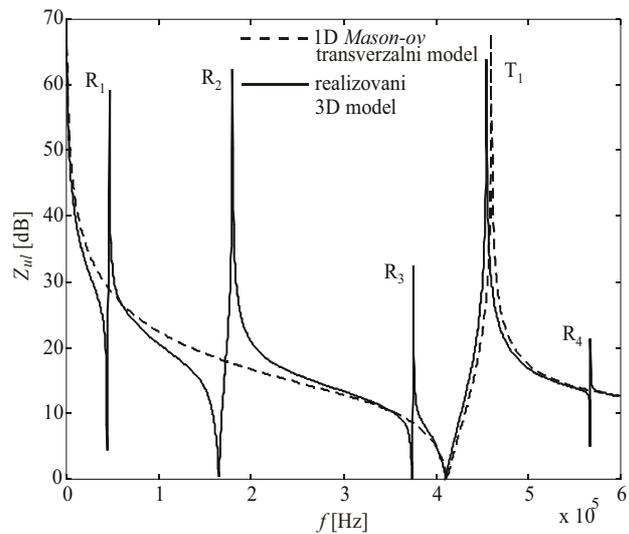


Figure 8.277. - Electric impedance $Z_{ul} = Z_{ul}(f)$

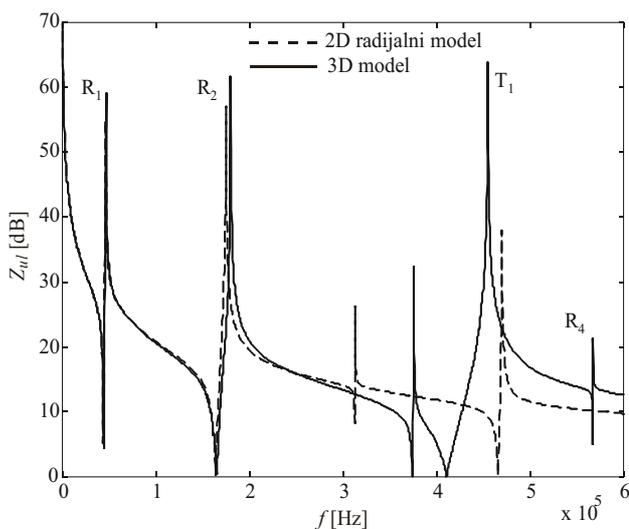


Figure 8.278. - Electric impedance $Z_{ul} = Z_{ul}(f)$

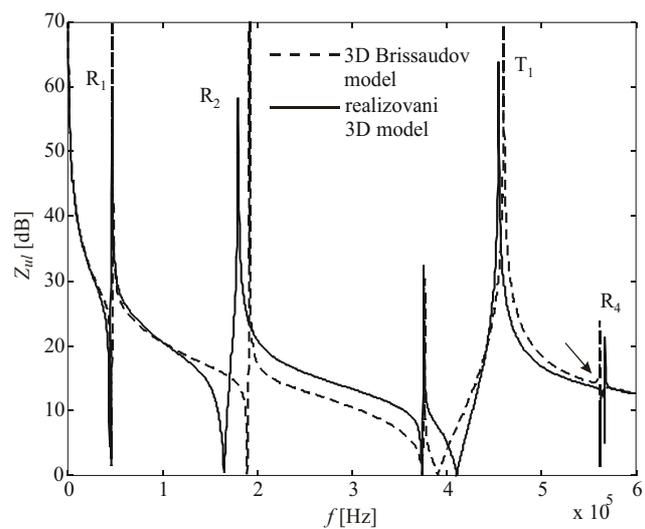


Figure 8.279. - Electric impedance $Z_{ul} = Z_{ul}(f)$

* 8.5.4. Second approach to the problem of 3D oscillations of circularring plate.

On Figure 8.280. is presented *characteristic of electric impedance dependence of frequency* for the same PZT8 piezoceramic specimen, thereat now are contour surfaces of the circular-ring plate loaded with different external loads (different acoustic impedances). Solid line represents case of piezoceramic circular-ring plate that oscillates freely in air without additional external load. Dashed line represents case when top metalized surface of the circular-ring plate ($z = h$) is loaded, while other surfaces are free. Dotted line is case of simultaneous loading of both top and bottom metalized surface $z = \pm h$. From Figure 8.280 one may see that *acoustic load* of the circular-ring plate in transversal direction affects mostly the transversal (crossing) mode of oscillation, while its influence on radial modes may be neglected.

On Figure 8.281. is presented case of loading of cylindric contour surfaces of circular-ring plate in radial direction. One may notice that increase of acoustic load in radial direction considerably affects radial resonant modes of oscillation of circular-ring plate, while influence on transversal modes may be neglected. Solid line represents *circular-ring plate* that oscillates freely in air without additional external load. Dashed line represents action of acoustic load on internal cylindric surface $r = a_1$, and dotted line simultaneous action of acoustic load on internal and external cylindric surface ($r = a_1$ and $r = a_2$).

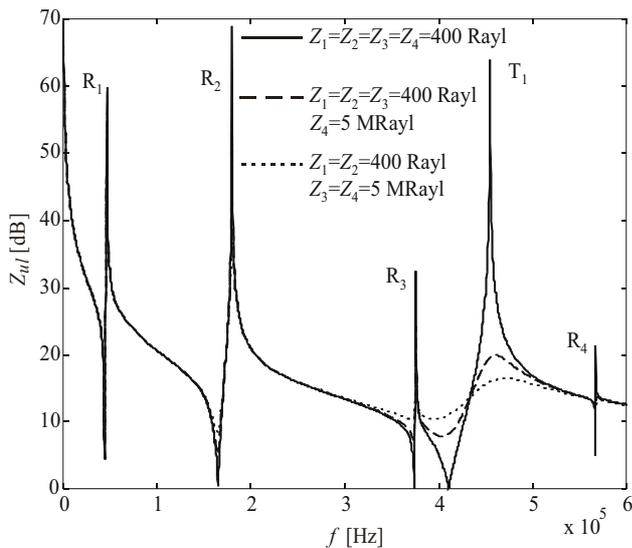


Figure 8.280. - Electric impedance $Z_{ul} = Z_{ul}(f)$

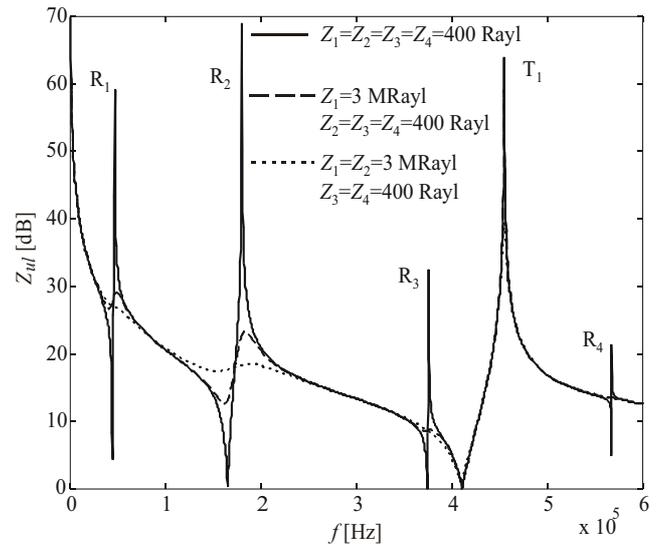


Figure 8.281. - Electric impedance $Z_{ul} = Z_{ul}(f)$

In order to further represent capabilities of proposed 3D model on Figure 8.282. and 8.283., *dependence of electric impedance in function of frequency f and thickness of the ring $2h$* is presented.

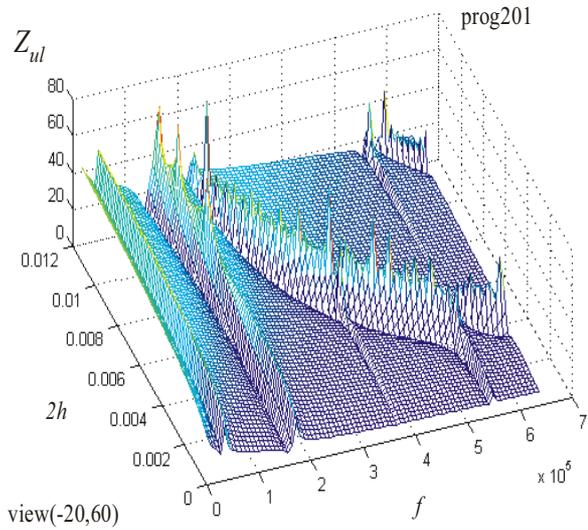


Figure 8.282. - Electric impedance $Z_{ul} = Z_{ul}(f, 2h)$

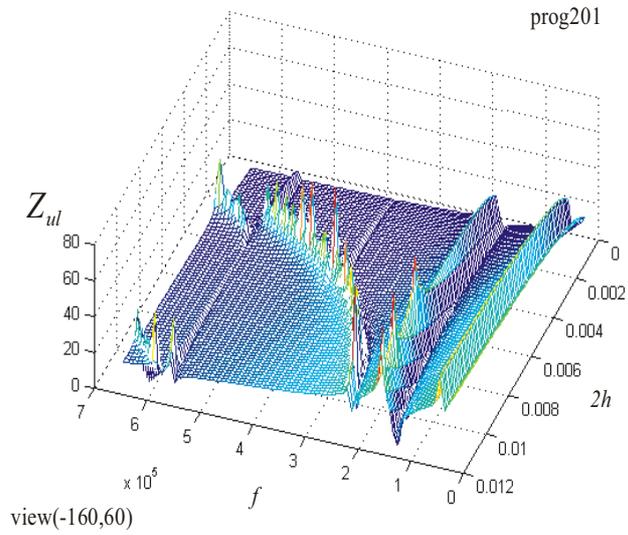


Figure 8.283. - Electric impedance $Z_{ul} = Z_{ul}(f, 2h)$

It is adopted that thickness of the circular-ring plate is in range $2h = 0 \div 12 \text{ mm}$. From Figure 8.282. one may notice that increase of the ring thickness affects mostly the *transversal (crossing) resonant mode* of oscillation, which shifts to the region of lower frequencies, and less the radial resonant normal modes of oscillation. *Radial resonant modes* shift too, but it is a result of mutual coupling with transversal resonant mode of oscillation. Also, change of thickness of circular-ring plate affects *value of capacitance of piezoceramic ring*, which reflects on the change of level height of electric impedance (Figure 8.282 and Figure 8.283).

On Figure 8.284. and 8.285. is presented *input electric impedance* Z_{ul} in function of frequency f and ratio of internal and external radius a_1 / a_2 (range is $a_1 / a_2 = 0 \div 1$).

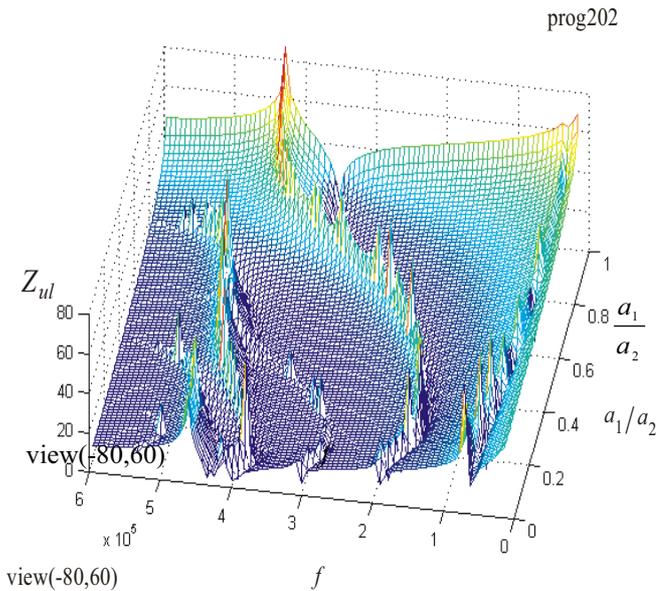


Figure 8.284. - Electric impedance $Z_{ul} = Z_{ul}(f, a_1 / a_2)$

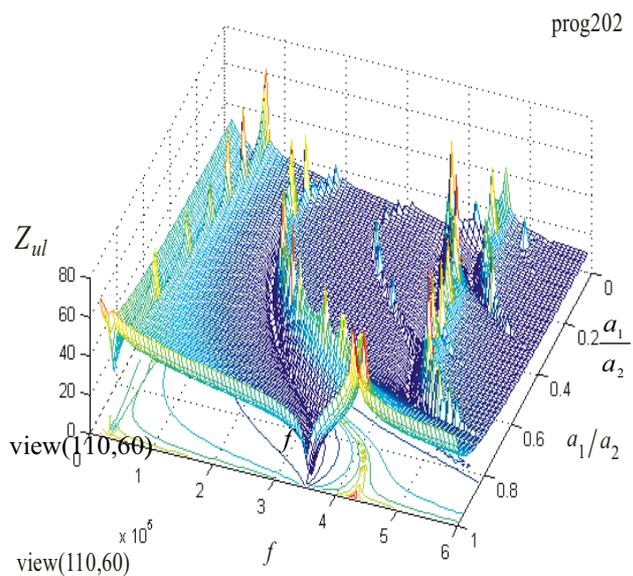


Figure 8.285. - Electric impedance $Z_{ul} = Z_{ul}(f, a_1 / a_2)$

As expected, one may see (Figure 8.284. and Figure 8.285) that change of radius ratio has the greatest influence on state of radial resonant modes, thereat one may notice an interesting phenomenon that first *radial resonant mode* moves to lower frequencies, while other radial resonant modes tend to higher frequencies, and thereat considerably affect the *transversal resonant mode*. This influence of internal radius on resonant frequency of transversal oscillations was not analyzed till now in the field of modelling of piezoceramic circular-ring plates and ultrasonic sandwich transducers. Similarly to the previous case, alteration of area size of metalized surface due to the alteration of internal radius generates alteration of its capacitance, and by itself also change of level value of input electric impedance.

Results from previous analysis can be presented more clearly through diagrams of spatial states of input electric impedance Z_{ul} in function of frequency and applied external load (applied acoustic impedance) Z_3 and Z_4 . Spatial diagrams of input electric impedance presented on Figure 8.286. and 8.287 conform to the input electric impedance from Figure 8.280., while spatial diagrams of input electric impedance presented on Figure 8.288. and 8.289 conform to the input electric impedance from Figure 8.281. On these Figures are clearly noticed values of external (acoustic) loads at which some *resonant modes* disappear.

Transversal (crossing) mode of oscillation disappears with increase of external load in direction of polarization axis across thickness of circular-ring plate (Figure 8.286. and Figure 8.287). Radial mode of oscillation disappears with increase of acoustic load Z_1 and Z_2 in radial direction on cylindric surfaces of circular-ring plate (Figure 8.288. and Figure 8.289). However, since it is about coupled tensors of state, these influences are not isolated (independent), but also the *resonant modes are coupled*. From Figure 8.286. and 8.287. One may see that changes at transversal resonant mode also affect changes of the third and fourth radial mode of oscillation, which are closest to the transversal mode, and at high external loads they also affect changes of distant radial resonant mode of oscillation.

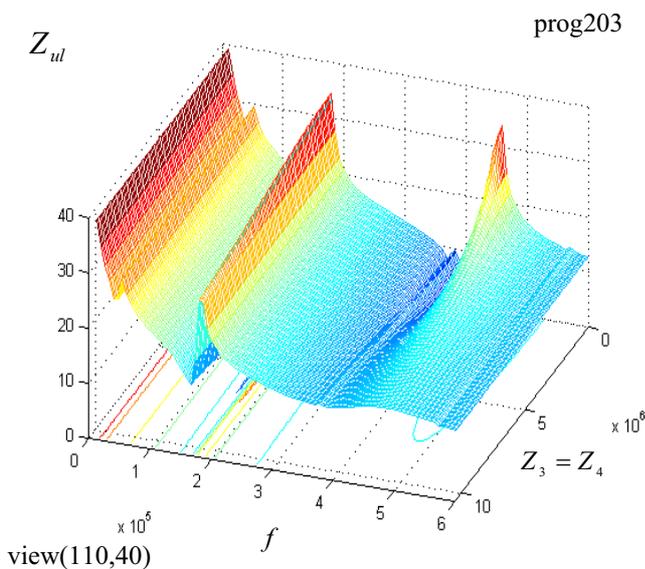


Fig. 8.286. - Electric impedance $Z_{ul} = Z_{ul}(f, Z_3 = Z_4)$

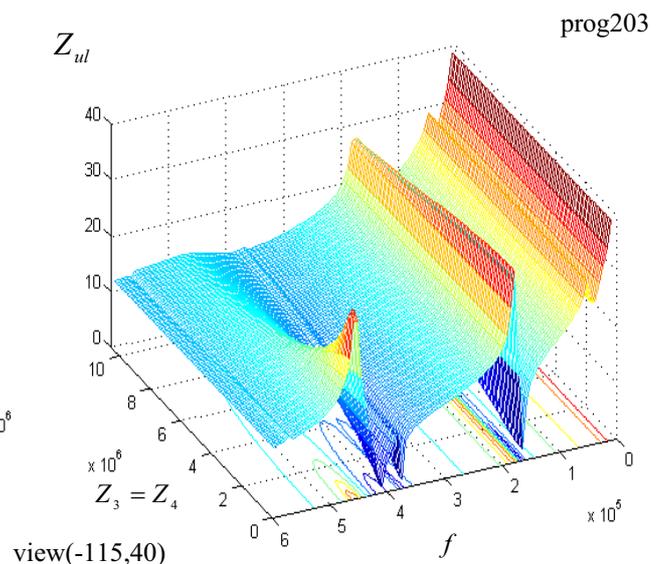


Fig. 8.287. - Electric impedance $Z_{ul} = Z_{ul}(f, Z_3 = Z_4)$

Z_{ul} prog204

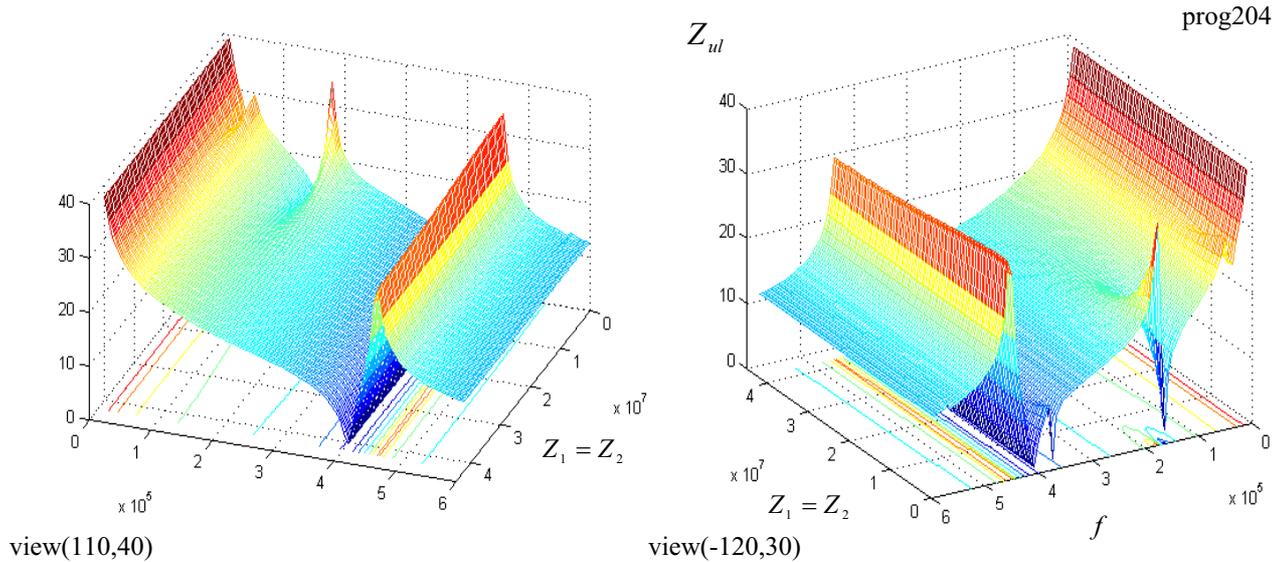


Fig. 8.288. - Electric impedance $Z_{ul} = Z_{ul}(f, Z_1 = Z_2)$ Fig. 8.289. - Electric impedance $Z_{ul} = Z_{ul}(f, Z_1 = Z_2)$

Mutual arrangement of radial and transversal resonant modes of oscillation depend of relation between internal and external radius, as well as of thickness of the piezoceramic circular-ring plate itself (figures 8.282, 8.283, 8.284, and 8.285). For definite concrete dimensions of piezoceramic circular-ring plate *resonant modes of oscillation* may be very close, and analyzed mechanical external load caused by acoustic impedances, can generate, besides intensity decrease of resonant modes, also their *frequency shift*. So, for example, at circular-ring plate with small external diameter and greater height, cylindric lateral surfaces become dominant, and lateral external radial loads (acoustic impedances) have significant influence on radial resonant modes of oscillation. In these cases, disappearing of radial modes with increase of external load can generate *frequency shift* of transversal mode of oscillation because of their very significant mutual coupling (*coupled tensors of piezoelectric material state*).

8.5.5.2. Analysis of Numerical and Experimental Results

Main task was to obtain experimental verification of adopted and analyzed 3D model* for coupled tensors of state of piezoelectric materials, what was the ultimate goal of this dissertation. *Input electric impedance* in function of frequency $Z_{ul} = Z_{ul}(f)$ is measured for different piezoceramic elements in shape of circular-ring plates and circular plates. Obtained experimental results are compared with correspondent results for adopted model obtained by using of software package MATLAB (Figure 8.290). In experimental and numerical analysis are used two types of piezoceramic materials, PZT4 and PZT8, whose *tensors of material coefficients* are presented in Appendix II. Here are observed cases of oscillation of piezoceramic specimens by excitation, i.e., by leading of electric energy from alternate voltage generator on electrode coatings that are located on principal mutually parallel plane surfaces, and which are perpendicular to the axis of polarization (Figure 8.121). In Appendix III are presented *characteristic programmes of numerical analysis using software package MATLAB for proposed and analyzed model* of circular-ring plate from Figure 8.121, and that is an *universal 3D model*, because it very well numerically simulates mutually coupled tensors at piezoceramic circular-ring plates and circular plates, as in plane,

* 8.5.4. Second approach to the problem of 3D oscillations of circular-ring plate.

so in space. Other programmes used in this dissertation are similar, and they are not presented because of volume and complexity of exposed matter. Dependence of electric impedance of frequency, for piezoceramic circular-ring plates and circular plates, is experimentally measured by *automatic network analyzer* HP4194A*, and comparative results are presented on Figures 8.291, 8.293, 8.295, 8.297, 8.299, 8.301.

On Figure 8.290. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer* in function of frequency, for PZT8 piezoceramic circular-ring plate with dimensions $2a_1 = 4$ [mm], $2a_2 = 10$ [mm], and $2h = 2$ [mm]. It is assumed that oscillation is performed in air without additional external loading.

On Figure 8.291. is presented comparison of quoted numerically simulated and measured experimental dependence of modulus of input electric impedance.

From Figure 8.290. and 8.291. one may see that forms and calculated values of input electric impedances, as well as calculated values of resonant and antiresonant frequency of oscillation for circular-ring piezoceramic plate, are very close to the correspondent experimentally obtained results, as for the *first radial mode of oscillation* R_1 , as well as for the *first transversal (crossing) mode of oscillation* T_1 .

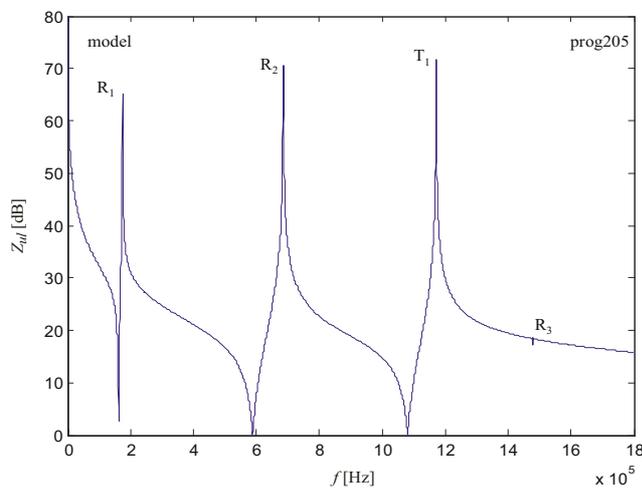


Figure 8.290. - Electric impedance $Z_{ul} = Z_{ul}(f)$

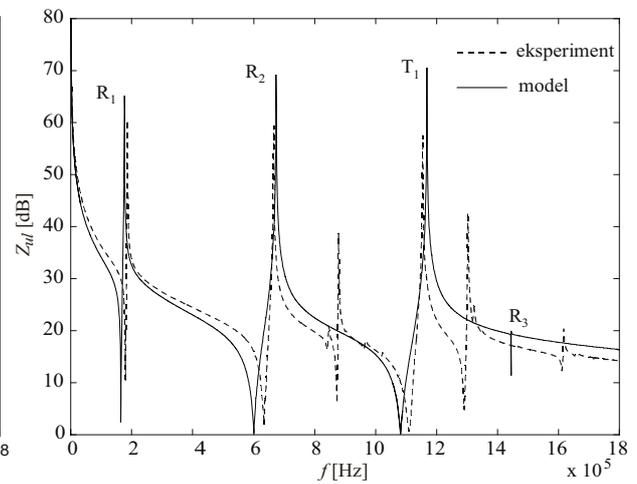


Figure 8.291. - Electric impedance $Z_{ul} = Z_{ul}(f)$

On Figure 8.292. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer* in function of frequency, for PZT4 piezoceramic circular-ring plate with dimensions $2a_1 = 13$ [mm], $2a_2 = 38$ [mm], and $2h = 4$ [mm]. It is assumed that oscillation is performed in air without additional external loading.

On Figure 8.293. is presented comparison of *characteristic of modulus of input electric impedance numerically simulated on computer, and experimental characteristic measured on automatic network analyzer* HP4194A.

* Network Impedance Analyzer.

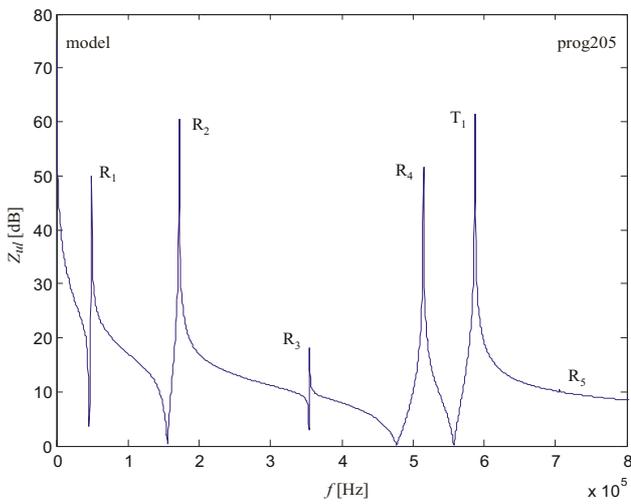


Figure 8.292. - Electric impedance $Z_{ul} = Z_{ul}(f)$

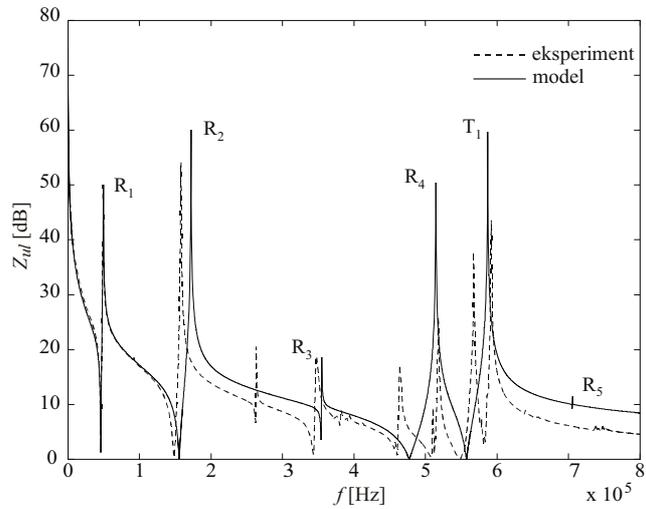


Figure 8.293. - Electric impedance $Z_{ul} = Z_{ul}(f)$

On Figure 8.294. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer* in function of frequency, for PZT4 piezoceramic circular-ring plate with dimensions $2a_1 = 13$ [mm], $2a_2 = 38$ [mm], and $2h = 6,35$ [mm]. It is assumed that oscillation is performed in air without additional external loading.

On Figure 8.295 is presented comparison of simulated and experimentally measured dependence of modulus of input electric impedance.

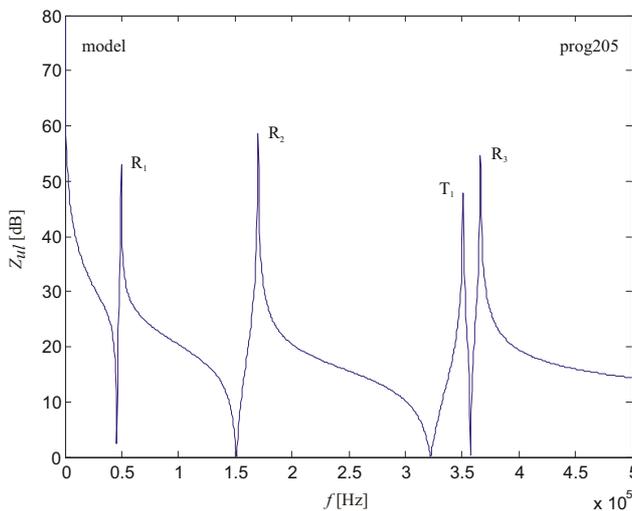


Figure 8.294. - Electric impedance $Z_{ul} = Z_{ul}(f)$

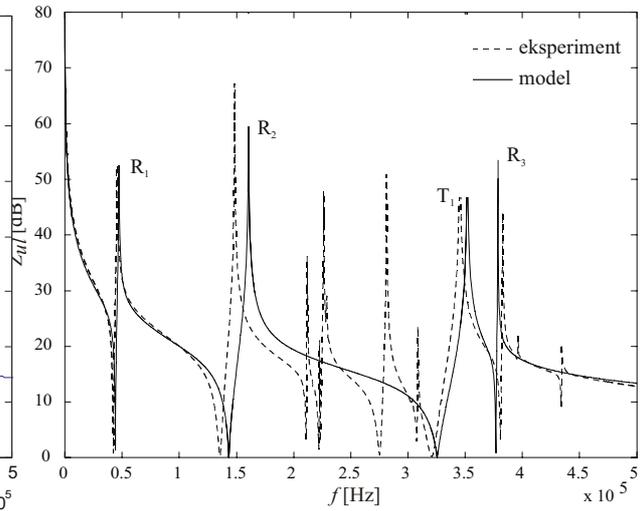


Figure 8.295. - Electric impedance $Z_{ul} = Z_{ul}(f)$

First radial mode R_1 and *first transversal mode* T_1 are the most often-used modes of oscillation at piezoelectric transducers in practical application. Adopted and analyzed model, as one may see, predicts with very good accuracy these modes of oscillation at different types of piezoceramic circular-ring specimens. *Transversal mode of oscillation* T_1 is most frequently used at ultrasonic high frequency

transducers. However, there is often need in practice for use of transducers at lower frequencies. Obtaining of lower operating resonant frequencies is possible to realize by application of *Langevin's* sandwich transducer, or simple piezoceramic circular-ring plate (or circular plate), which oscillates in its first radial mode R_1 . At first radial resonant mode R_1 a considerable stressing in transversal (crossing) direction exists because of the elastic coupling and high interaction between coupled *tensors of state of piezoelectric materials*. Analysis of the cited transversal oscillation of circular-ring piezoceramic plate is enabled by proposed 3D model, and it may be broadened on cylindric piezoceramic bodies in shape of circular plates. Adopted 3D model considers mutual coupling of transversal and radial oscillations, and thereby is possible to determine optimal geometry of circular-ring plate or circular plate, in order to obtain increased displacement in transversal direction during oscillation in region of the first radial mode.

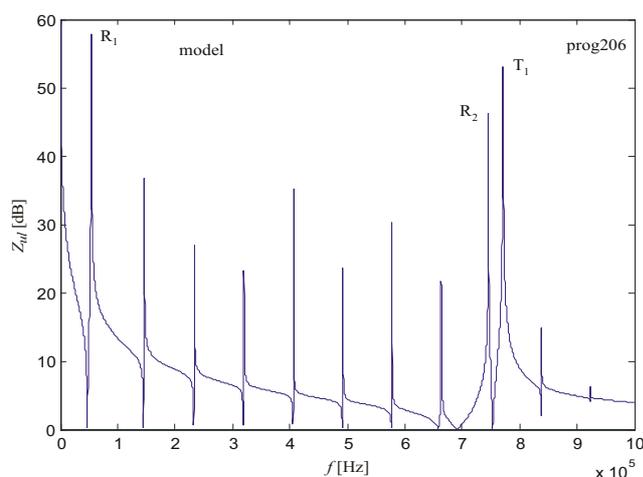


Figure 8.296. - Electric impedance $Z_{ul} = Z_{ul}(f)$

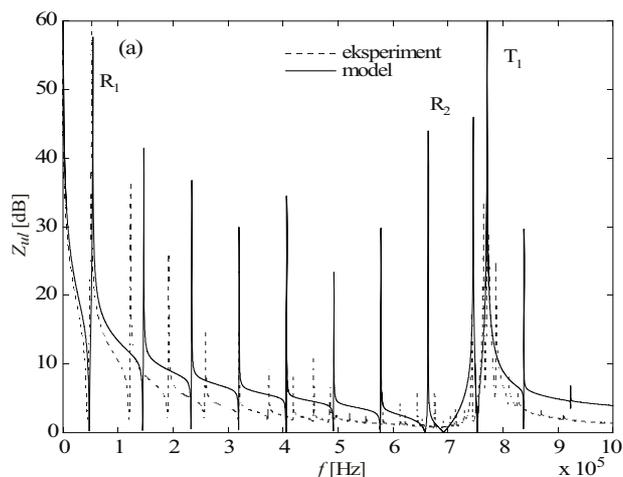
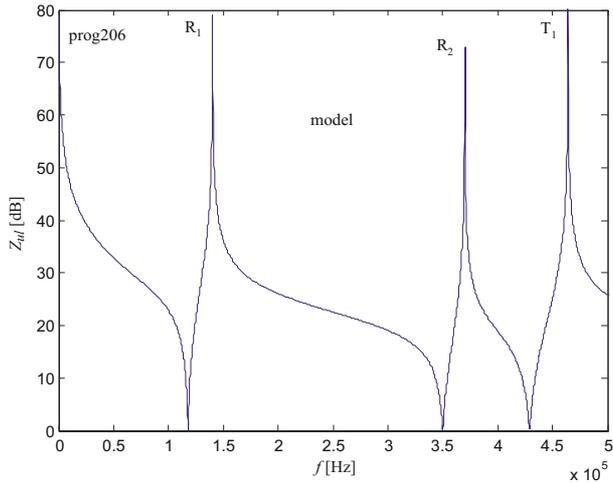
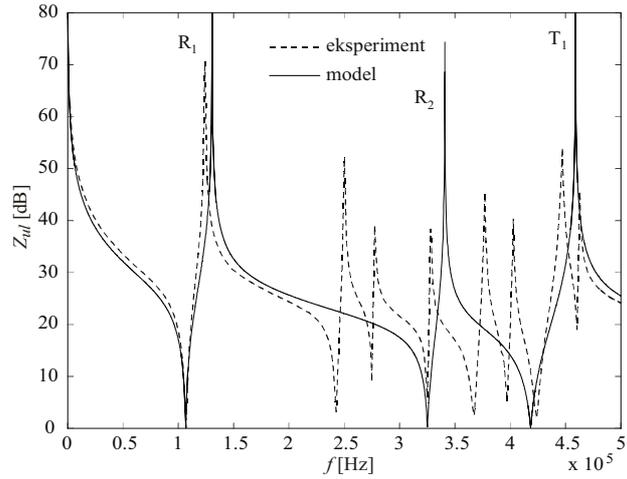


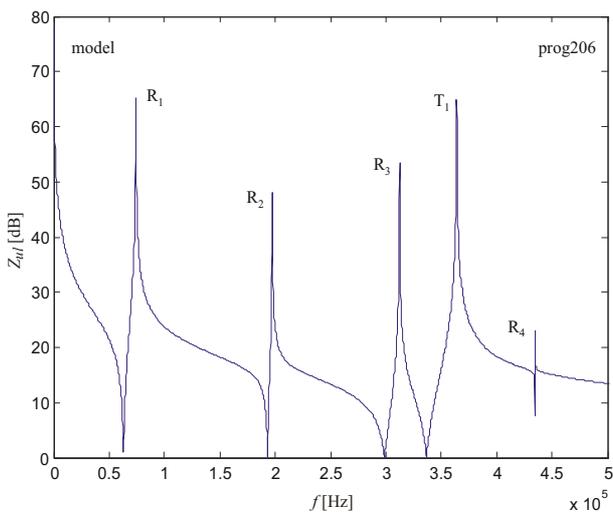
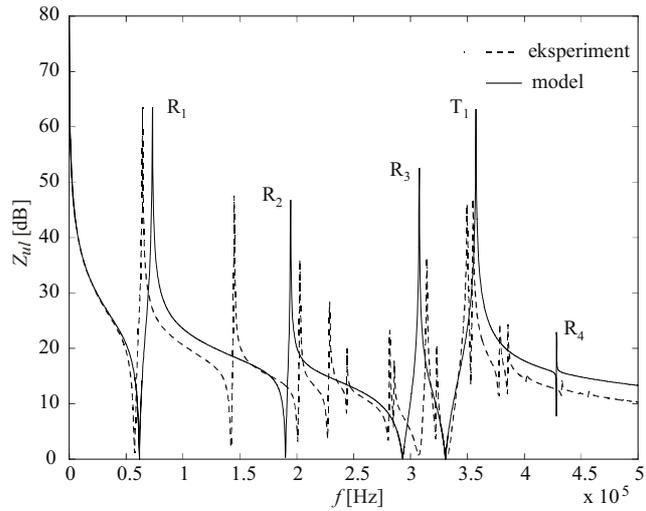
Figure 8.297. - Electric impedance $Z_{ul} = Z_{ul}(f)$

On Figure 8.296. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer* in function of frequency, for PZT8 piezoceramic circular plate with dimensions $2a_2 = 50$ [mm], and $2h = 3$ [mm]. It is assumed that oscillation is performed in air without additional external loading. On Figure 8.297. is presented comparison of numerically simulated and experimentally measured dependence of modulus of input electric impedance.

On Figure 8.298. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer* in function of frequency, for PZT4 piezoceramic circular plate with dimensions $2a_2 = 20$ [mm], i $2h = 5$ [mm]. It is assumed that oscillation is performed in air without additional external loading. On Figure 8.299. is presented comparison of numerically simulated and experimentally measured dependence of modulus of input electric impedance.

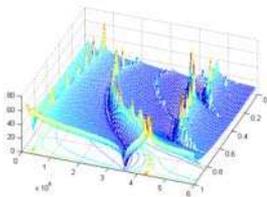
Figure 8.298. - Electric impedance $Z_{ul} = Z_{ul}(f)$ Figure 8.299. - Electric impedance $Z_{ul} = Z_{ul}(f)$

On Figure 8.300. is presented *characteristic of modulus of input electric impedance, simulated numerically on computer in function of frequency, for PZT4 piezoceramic circular plate with dimensions $2a_2 = 38$ [mm], and $2h = 6,35$ [mm]. It is assumed that oscillation is performed in air without additional external loading. On Figure 8.301. is presented comparison of numerically simulated on computer and measured experimental dependence of modulus of input electric impedance.*

Figure 8.300. - Electric impedance $Z_{ul} = Z_{ul}(f)$ Figure 8.301. - Electric impedance $Z_{ul} = Z_{ul}(f)$

In case of piezoceramic circular plates, experimentally measured and numerically modelled dependences of electric impedance of frequency, in range of the first transversal mode T_1 , even better coincide regarding the cases of circular-ring specimens, and satisfying results are also achieved in range of the first radial mode R_1 . It is noticed that, considering *higher radial resonant modes* of oscillation $R_2, R_3, R_4, R_5, \dots$, *modelled resonant frequencies* of radial modes are mostly greater than measured ones, and rarely smaller than measured resonant frequencies (this conclusion stands for circular-ring plates and circular piezoceramic plates). One of possible causes for arising of this phenomenon is real presence of other types of oscillatory resonant modes, which are not encompassed by proposed and adopted model (for example: *edge mode of oscillation* and *flexion mode of oscillation – thickness-shear*), and which occur in region between the first radial R_1 and first transversal T_1 resonant mode of oscillation. Presence of other types of oscillatory modes is especially characteristic at piezoceramic specimens in shape of circular plates. This observation is best seen on Figure 8.299, where some resonant modes on experimental characteristic, in the vicinity of the modelled radial mode R_2 , do not represent radial resonant

modes. Because the adopted model did not encompass these types of oscillation, *numerically modelled dependence of impedance* was mostly above experimentally measured characteristic, and it did not descend even at higher frequencies. Exception is the case of piezoceramic circular plate from Figure 8.297., where missing and not encompassed by model modes are poorly coupled with oscillatory modes present in model, so they have not great influence on electric impedance characteristic. Second possible cause of arising of the mentioned phenomenon is that model does not consider local mechanical and dielectric losses, occurrence of heating, and electrostriction of the piezoceramic element. Minimum and maximum values of electric impedance at resonant frequencies of radial and transversal oscillation are more distinct at calculated numerically modelled characteristic, regarding the experimentally obtained characteristic measured on automatic network analyzer HP4194A, at all analyzed cases of piezoceramic circular-ring plates and circular plates. Nevertheless all cited in the analysis, one can make a positive conclusion that, by proposed *and adopted 3D model*, generally in advance, even at design stage, one may predict the state of electric impedance with precise determination of frequencies of dominant, mutually coupled, resonant oscillatory modes. This conclusion is very important, as for theory and theory of experiment, as well as for the manufacturing technology, because one may predict behaviour of piezoceramic elements before their immediate workmanship, even at stage of calculation and design process.



Dr. Ljubisa Peric,

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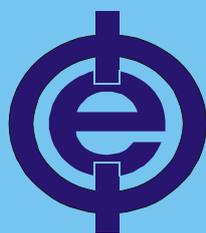
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University of Niš
Faculty of Electronics

Milan Đ. Radmanović, Dragan D. Mančić

**DESIGNING AND MODELLING OF
THE POWER ULTRASONIC
TRANSDUCERS**



Edition: Monographies



A SERIES OF EXTRAORDINARY AND UNIQUE BOOKS RECOMMENDED BY MPI

Dr. Milan Đ. Radmanović, Dr. Dragan D. Mančić

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PREFACE

Field of power ultrasonic technique, which represents an important field of industrial electronics, in recent two decades experienced very swift and dynamic development. An intensive development concerns as design and construction of new ultrasonic devices, as well as broadening of application fields of power ultrasound in many industrial branches and processes (mechanical, electric, and chemical industry). Aside with appearing of new applications of ultrasound, new, more perfect sandwich transducers are designed and developed, and numerous scientific papers appeared, in which are treated different aspects of power ultrasonic technique, especially different electromechanical models by which is obtained design and optimization of ultrasonic transducers.

In this monograph firstly is performed systematization of different existing procedures and methods for modeling of power ultrasonic transducers. Besides that, new procedures of modeling, design, and optimization of power ultrasonic transducers are presented, based on previously realized original models of piezoceramic and metal rings. Thus is completed design of a sandwich transducer as a unique system, consisted of piezoceramic rings, emitting and reflecting metal ending, as well as of central bolt. Basic idea of the authors was to help with realized models to the designers of new ultrasonic systems, due to the fact that currently there is no literature from this field in Serbian.

Original results, presented in this monograph, are product of several-year-research in the field of power ultrasound in the Laboratory for energetic electronics and control of electroenergetic transducers in the Faculty of Electronics in Niš, wherefrom originated over 50 scientific papers from this field. Concrete results, presented here, are part of one master thesis and one doctoral dissertation, realized in the frame of research in this field.

On this occasion authors express their gratitude to the reviewers, Prof Vanča Litovski, Ph.D. and Prof Stojan Ristić, Ph.D. on their useful suggestions and notes.

Niš, January 2004

Authors

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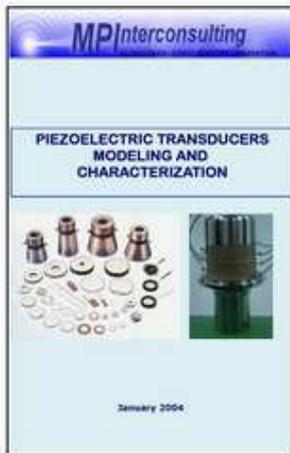
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PIEZOELECTRIC TRANSDUCERS **MODELING AND CHARACTERIZATION**

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1. Piezoelectric Converters Modeling and Characterization

Here are pictures of products that are related to modeling and characterization techniques presented in this paper. All of such devices and components can be measured, modeled and characterized using (more or less) the same methodology and the same terminology.



There are many possibilities to present and analyze equivalent models of piezoelectric converters. The analysis presented here will primarily cover modeling of high-power piezoelectric, Langevin, sandwich-converters (applicable in ultrasonic cleaning, plastic welding, sonochemistry and other power industrial applications), and the same analysis can easily be extended to cover much wider field of different piezoelectric transducers and sensors (but this was not the main objective of this paper). For electrical engineering needs (as for instance: when optimizing ultrasonic power supplies, in order to deliver maximal ultrasonic power to a mechanical load) we need sufficiently simple and practical (lumped parameter), equivalent models, expressed only using electrical (and easy measurable or quantifiable) parameters (like resistance, capacitances, inductances, voltages and currents). Of course, in such models we should (at least) qualitatively know which particular components are representing purely electrical nature of the converter, and which components are representing mechanical or acoustical nature of the converter, as well as to know how to represent mechanical load. For here described purpose, the best lumped parameter equivalent circuits that are fitting a typical piezoelectric-converter impedance (the couple of series and parallel resonance of an isolated vibration mode) are shown to be Butterworth-Van Dyke (**BVD**) and/or its electrical dual-circuit developed by Redwood (both of them derived by simplifying the Mason equivalent circuit and/or making the best piezoelectric impedance modeling based on experimental results and electromechanical analogies). In this paper the two of mutually equivalent (above mentioned, **and later slightly modified**), dual electrical models will be used to present a piezoelectric converter operating in its series and/or parallel resonance, Fig. 1.

Above described objective has been extremely simplified after electronics industry developed Network Impedance (Gain-Phase) Analyzers (such as HP 4194A and similar instruments). Practically, for the purpose of modeling, it is necessary to select one single converter's operating mode (to select a frequency window which captures only the mode of interest, or the single couple of series and parallel resonance belonging to that mode) and let Impedance Analyzer to perform electrical impedance measurements by producing sweeping frequency signal in the selected frequency interval (and by measuring voltage and current passing on the converter connected to the input of Impedance Analyzer). The next step (implemented in Impedance Analyzers) is to compare the measured impedance parameters with theoretically known converter model (lumped parameters model, already programmed as an modeling option inside of Impedance Analyzer), and to calculate model parameters (practically performing the best curve fitting that places measured impedance values into theoretical impedance model). This way, in a few (button pressing) steps we are able to get numerical values of all (R, L, C) electrical components relevant (only) for selected converter mode and selected frequency range (and this is in most cases the most important for different engineering purposes, such as: optimizing ultrasonic power supplies, realizing optimal resonant frequency and output power control, optimizing converters quality...). We are also able to compare (using modern impedance analyzers) how close are measured impedance values (of a real converter), and values resulting from impedance curve fitting process. In cases of well-designed converters (and converters with sufficiently high mechanical quality factor) we are able to get almost 100% correct modeling in a selected frequency window (meaning that all measured and calculated R, L and C, lumped model parameters, are numerically almost 100% correct).

The objectives of this paper are:

1. To explain the most important (and simplified, practical and easy quantifiable), electrical lumped-parameters equivalent circuits, suitable to represent piezoelectric converter in its series and/or parallel resonance, for purposes such as converters characterization, qualification and optimization, for different electrical design purposes, as well as to explain qualitatively converter models regarding higher frequency harmonics, and

models when converter transforms mechanical input into electrical output (operating as a receiver or sensor).

2. To establish the very general concept of mechanical loading of piezoelectric converters (where mechanical load is presented in normalized form using the mechanical-load units comparable to internal resistance of the converter-driving electric circuit, or ultrasonic power supply).
3. To analyze the optimal power transfer of piezoelectric converters (when transforming electrical input into mechanical output), operating in series and parallel resonance, and to explain mechanical loading process and losses in both situations.

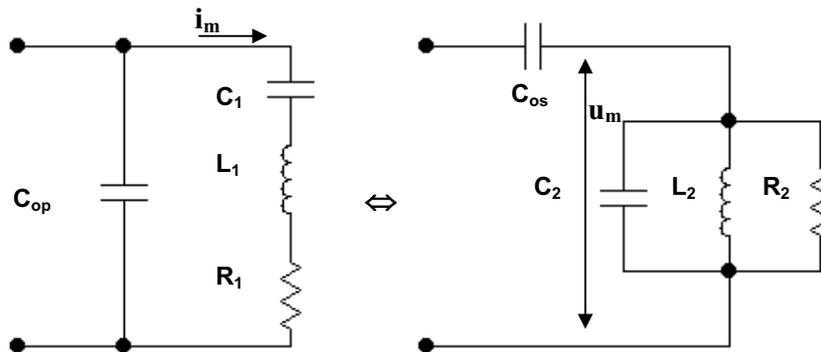


Fig. 1 Piezoelectric Converter (One-Port), Dual, BVD Models valid in the close vicinity of an isolated couple of converters series and parallel resonances

The Fig. 1 presents two of the most widely used lumped-parameters piezoelectric converter impedance models (mutually equivalent, **BVD = Butterworth-Van Dyke**, dual-circuits models), valid for isolated couple of series and parallel resonances (of a non-loaded). In fact, on the fig Fig. 1 are presented the simplest models applicable for relatively high mechanical quality factor piezoelectric converters, where thermal dissipative elements in piezoceramics could be neglected. The more general models (again mutually equivalent), representing real piezoelectric (non-loaded) converters with dissipative dielectric losses and internal resistive electrode-elements (R_{op} (=) **Leakage AC and DC resistance**, $R_{os1} \approx R_{os2}$ (=) **Dielectric resistive loss of piezoceramics**) in piezoceramics are presented on the Fig. 2. For high quality piezoceramics R_{op} , is in the range of 10 M Ω - 50 M Ω , and R_{os1} , R_{os2} are in the range between 50 Ω and 100 Ω , measured at 1 kHz, low signal (and can be calculated from piezoceramics $\tan\delta$ value, or using HP 4194A, and similar Impedance Analyzers). In most of cases of high quality piezoceramics we can neglect R_{op} as too high resistance, and R_{os1} , R_{os2} , as too low values, but we should also know that dielectric and resistive losses are becoming several times higher when converter is driven high power, in series or parallel resonance, comparing them to low signal measurements.

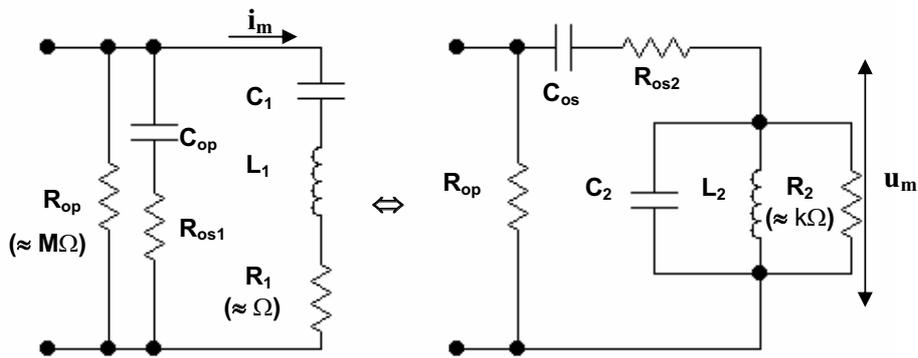


Fig. 2 BVD Piezoelectric Converter Models with dissipative elements

The other dissipative power losses (R_1 and R_2) are belonging to the mechanical circuit branch and come from converter joint losses, from planar friction losses between piezoceramics and metal parts, from mounting elements and from material hysteresis-related losses (internal mechanical damping in all converter parts). The models from the Fig. 2 can be schematically simplified if we introduce abbreviated electric-elements symbolic presenting dissipative (real) inductances and capacitances together with their belonging resistances, using only one symbol, as for instance: For any electric combination (or connection) between one capacitance and one resistance we shall introduce the symbol C^* , and for any electric combination between one inductance and one resistance we shall introduce the symbol L^* (since we can always find exact circuit transformations between two elements in serial and parallel connection). Doing this way, models presented on Fig. 2 will be simplified as given on the Fig. 3, and applicable circuit equivalents (used in Fig. 3) are presented on the Fig. 4.

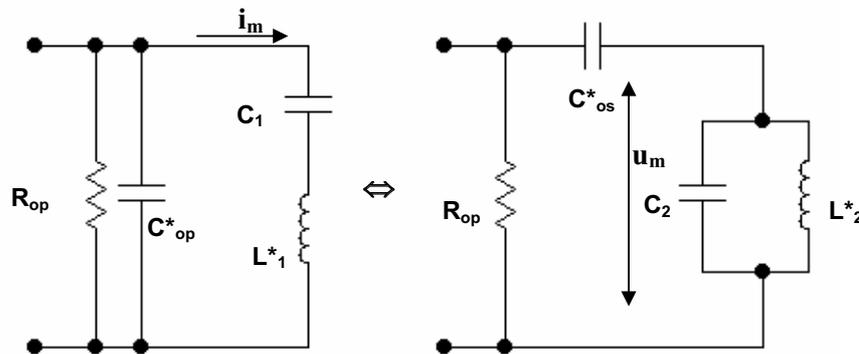


Fig. 3 Simplified BVD Piezoelectric Converter Models
 (L^* and C^* are presenting real inductances and capacitances with internally integrated, dissipative elements: Fully equivalent to models on Fig. 2)

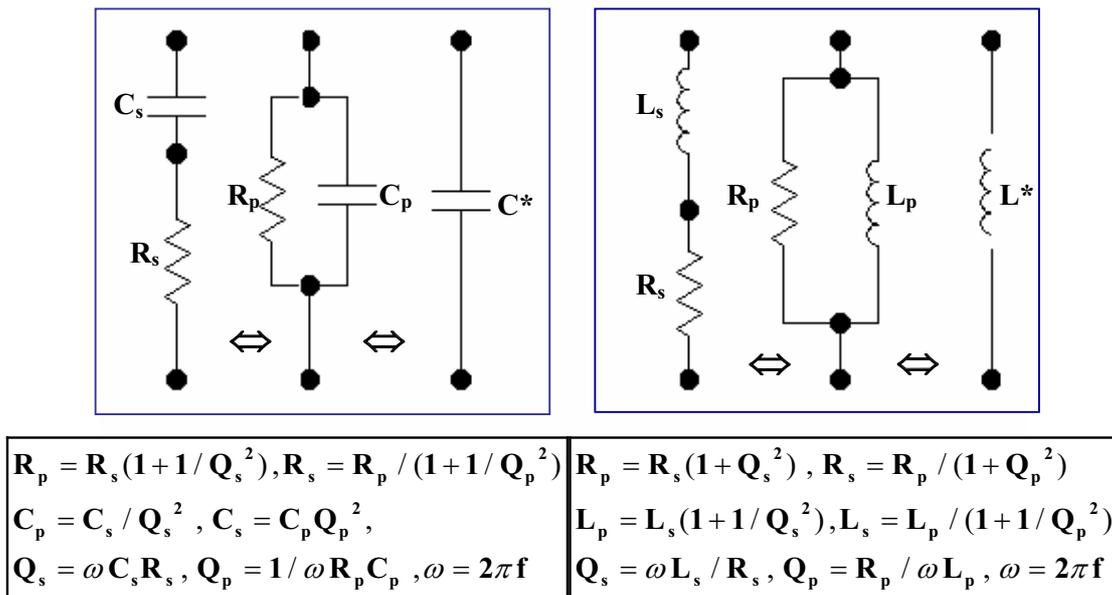


Fig. 4 Circuit Equivalents & Simplifications (explaining models from Fig. 3)

The other elements on the Fig. 2 are: $C_{os} \approx C_{op}$ (\Rightarrow) **Clamped, static capacitance/s of piezoceramics**, $C_{1/2}$, $L_{1/2}$ (\Rightarrow) **motional mass and stiffness elements of converter's mechanical oscillating circuit/s** (see Fig. 7 to find approximate mathematical relations between all model parameters). We could also add in series to any of input converter terminals the cable (and winding) resistance, since every real converter has input electrodes, soldered or bonded (electrical) joints, and a cable (presently neglected parameters).

The influence of an external acoustic load on the converters' modeling is presented on the Fig. 5, by introducing loading resistances R_{L1} and R_{L2} , as the closest and very much simplified equivalent of the real converter loading (in reality loading resistances R_{L1} and R_{L2} , sometimes should be treated as complex impedances as the most general case).

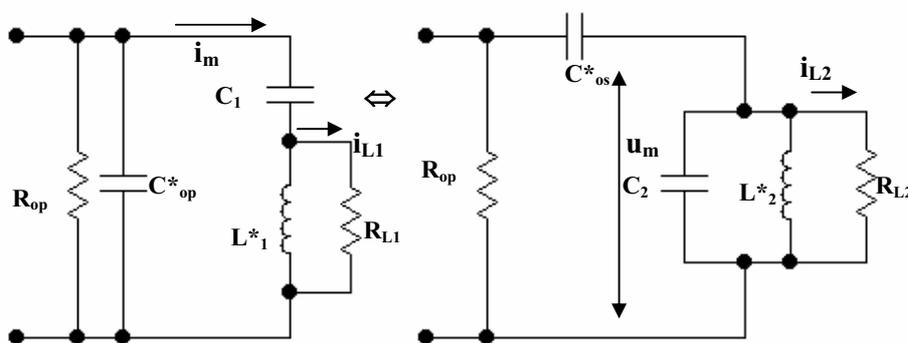


Fig. 5 Alternative BVD Models of Loaded Piezoelectric Converters

Based on equivalent electric circuits presented on Fig. 4, we can easily place parallel-loading resistances from Fig. 5 in series with inductances, just by calculating new equivalent frequency-dependant elements-values. In literature regarding the same problematic it is very usual to see that left-side piezoelectric converter-model from Fig. 5 has loading resistance in series with motional inductance and capacitance, and for the model on the right side of the Fig. 5 is usual that loading resistance is found in parallel with motional inductive and capacitive circuit elements (but using Electric Circuit Theory we can easily play with any of parallel or series elements combination, as presented on the Fig. 4). It is also clear that loading nature or load-resistance would change, depending how and where we place it (in situations, like in series connection/s with motional inductance, load resistance would increase with load-increase (starting from very low value), and in case of placing it in parallel with motional inductance (as presented on Fig. 5), load resistance would decrease with load-increase (starting from very high value)).

*In all above given converter models (Figs. 1,2,3,5), we can recognize motional current i_m and motional voltage u_m as the most important mechanical-output power/amplitude controlling parameters of piezoelectric converters in series and parallel resonance. When converter is operating in series resonance, in order to control its output power and/or amplitude we should control its motional current i_m , and in the regime of parallel resonance, output power and/or amplitude are directly proportional to the motional voltage u_m . More precisely, **when we compare two operating regimes of the same converter, when converter is producing the same output power (in series and/or parallel resonance)**, we can say that converter operating in series resonance is able to deliver to its load high output force (or high pressure) and relatively low velocity, and when operating in parallel resonance it is able to deliver high output velocity and relatively low force (knowing that output converter power is the product between velocity and force delivered on its front emitting surface). Here we are using the electromechanical analogy system: **(CURRENT \leftrightarrow FORCE) & (VOLTAGE \leftrightarrow VELOCITY)**. When we are talking about converter's series-resonance frequency zone, this is the case of motional **Current-Force** resonance (where converter's impedance has low values), and when we are talking about converter's parallel-resonance frequency zone, this is the case of motional **Voltage-Velocity** resonance (where converter has high impedance values). Automatically, if we realize by electrical means high motional current (current resonance, equal to series resonance), the converter will produce high motional force (it will operate in a force resonance). Also, if we realize by electrical means high motional voltage (voltage resonance, equal to parallel resonance), the converter will produce high motional velocity (it will operate in a velocity resonance). All above conclusions, for the time being, are based only on the analogy **(CURRENT \leftrightarrow FORCE) & (VOLTAGE \leftrightarrow VELOCITY)**, and later on, some more (experimental) supporting facts will be presented.*

It is also important to underline which circuit-elements (in all above found circuits, Figs. 1,2,3,5) are representing purely electrical elements of a piezoelectric converter, and which elements are only given as functional (and analog) electrical equivalents of converter's mechanical parts and its mechanical properties (including loading elements), see Fig.6.

It is very important to know that mechanical converter-loading, presented on Figs. 5 & 6, is equally and coincidentally influencing changes, both in series and parallel converter impedance, basically reducing equivalent mechanical quality factor/s of a loaded converter (or coincidentally increasing its series resonant-impedance and decreasing parallel resonant-impedance). ***This is the principal reason why (in this paper) an isolated couple of series and parallel resonances is treated as the same, single and unique oscillating-mode that can be driven in its current or voltage resonance, and produce force or velocity-dominant mechanical output. It is also shown to be possible to drive an ultrasonic converter high power, extremely efficiently, in any frequency (continuously) between its series and parallel resonance (when special converter reactive impedance-compensation is used).*** Since there is certain frequency shift between each couple of series and parallel resonances, and since most of today's converters (made by big players in ultrasonic industry worldwide) operate either in series or parallel resonance, in literature regarding converters modeling and converters measurements, many authors are talking about different vibration modes or different harmonics.